

# Eigenvalues of conical waveguides

**Konstantin Pankrashkin**

Laboratoire de mathématiques  
Université Paris-Sud, Orsay, France

Based on collaborations with  
Vincent Bruneau, Nicolas Popoff,  
Thomas Ourmières-Bonafos

# Infinite discrete spectrum

We are interested in operators whose discrete eigenvalues accumulate to the edges of the continuous spectrum.

# Infinite discrete spectrum

We are interested in operators whose discrete eigenvalues accumulate to the edges of the continuous spectrum.

- $A = -\Delta + V(x)$  in  $L^2(\mathbb{R}^d)$  with a long range potential,  $V(x) \sim -\frac{a}{|x|^{2-\varepsilon}}$  at  $\infty$ : one has the continuous spectrum  $[0, +\infty)$  and infinitely many negative eigenvalues.

**Example:** Coulomb potential  $V(x) = -1/|x|$  in  $L^2(\mathbb{R}^3)$ , the eigenvalues:  $-\frac{1}{4n^2}$ .

**Weyl asymptotics:** the number of eigenvalues below  $-\varepsilon$ ,  $N(A, -\varepsilon)$ , satisfies

$$N(A, -\varepsilon) \sim c_d \int_{\mathbb{R}^d} (V(x) - \varepsilon)_+^{d/2} dx, \quad \varepsilon \rightarrow 0.$$

# Infinite discrete spectrum

We are interested in operators whose discrete eigenvalues accumulate to the edges of the continuous spectrum.

- $A = -\Delta + V(x)$  in  $L^2(\mathbb{R}^d)$  with a long range potential,  $V(x) \sim -\frac{a}{|x|^{2-\varepsilon}}$  at  $\infty$ : one has the continuous spectrum  $[0, +\infty)$  and infinitely many negative eigenvalues.

**Example:** Coulomb potential  $V(x) = -1/|x|$  in  $L^2(\mathbb{R}^3)$ , the eigenvalues:  $-\frac{1}{4n^2}$ .

**Weyl asymptotics:** the number of eigenvalues below  $-\varepsilon$ ,  $N(A, -\varepsilon)$ , satisfies

$$N(A, -\varepsilon) \sim c_d \int_{\mathbb{R}^d} (V(x) - \varepsilon)_+^{d/2} dx, \quad \varepsilon \rightarrow 0.$$

- 1D model:  $T = -\frac{d^2}{dx^2} + V(x)$  in  $L^2(\mathbb{R}_+)$  with  $V \sim -\frac{c}{x^2}$  at  $\infty$ :  $[0, \infty) \subset$  spectrum,

$$N(T, -\varepsilon) \sim \frac{1}{2\pi} \sqrt{\left(c - \frac{1}{4}\right)_+} |\ln \varepsilon|, \quad \varepsilon \rightarrow 0.$$

# Infinite discrete spectrum

We are interested in operators whose discrete eigenvalues accumulate to the edges of the continuous spectrum.

- $A = -\Delta + V(x)$  in  $L^2(\mathbb{R}^d)$  with a long range potential,  $V(x) \sim -\frac{a}{|x|^{2-\varepsilon}}$  at  $\infty$ : one has the continuous spectrum  $[0, +\infty)$  and infinitely many negative eigenvalues.

**Example:** Coulomb potential  $V(x) = -1/|x|$  in  $L^2(\mathbb{R}^3)$ , the eigenvalues:  $-\frac{1}{4n^2}$ .

**Weyl asymptotics:** the number of eigenvalues below  $-\varepsilon$ ,  $N(A, -\varepsilon)$ , satisfies

$$N(A, -\varepsilon) \sim c_d \int_{\mathbb{R}^d} (V(x) - \varepsilon)_+^{d/2} dx, \quad \varepsilon \rightarrow 0.$$

- 1D model:  $T = -\frac{d^2}{dx^2} + V(x)$  in  $L^2(\mathbb{R}_+)$  with  $V \sim -\frac{c}{x^2}$  at  $\infty$ :  $[0, \infty) \subset$  spectrum,

$$N(T, -\varepsilon) \sim \frac{1}{2\pi} \sqrt{\left(c - \frac{1}{4}\right)_+} |\ln \varepsilon|, \quad \varepsilon \rightarrow 0.$$

- Geometric operators: waveguides  $\sim$  Dirichlet Laplacians in (long-range) perturbations of the infinite cylinders  $\Omega \times \mathbb{R}$  (Briet, Raikov, Soccorsi, ...).

# Infinite discrete spectrum

We are interested in operators whose discrete eigenvalues accumulate to the edges of the continuous spectrum.

- $A = -\Delta + V(x)$  in  $L^2(\mathbb{R}^d)$  with a long range potential,  $V(x) \sim -\frac{a}{|x|^{2-\varepsilon}}$  at  $\infty$ : one has the continuous spectrum  $[0, +\infty)$  and infinitely many negative eigenvalues.

**Example:** Coulomb potential  $V(x) = -1/|x|$  in  $L^2(\mathbb{R}^3)$ , the eigenvalues:  $-\frac{1}{4n^2}$ .

**Weyl asymptotics:** the number of eigenvalues below  $-\varepsilon$ ,  $N(A, -\varepsilon)$ , satisfies

$$N(A, -\varepsilon) \sim c_d \int_{\mathbb{R}^d} (V(x) - \varepsilon)_+^{d/2} dx, \quad \varepsilon \rightarrow 0.$$

- 1D model:  $T = -\frac{d^2}{dx^2} + V(x)$  in  $L^2(\mathbb{R}_+)$  with  $V \sim -\frac{c}{x^2}$  at  $\infty$ :  $[0, \infty) \subset$  spectrum,

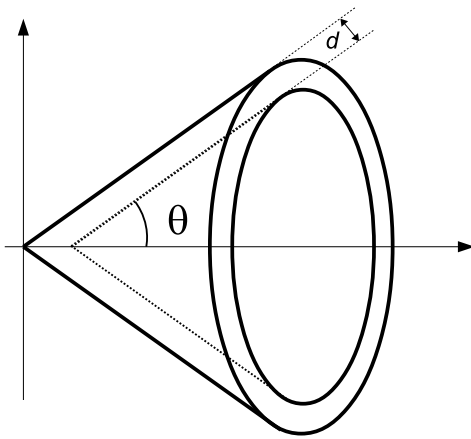
$$N(T, -\varepsilon) \sim \frac{1}{2\pi} \sqrt{\left(c - \frac{1}{4}\right)_+} |\ln \varepsilon|, \quad \varepsilon \rightarrow 0.$$

- Geometric operators: waveguides  $\sim$  Dirichlet Laplacians in (long-range) perturbations of the infinite cylinders  $\Omega \times \mathbb{R}$  (Briet, Raikov, Soccorsi, ...).

We are going to study the eigenvalues for several models involving “conical waveguides” whose cross-section grows linearly at infinity.

# Previous works

Conical layer:

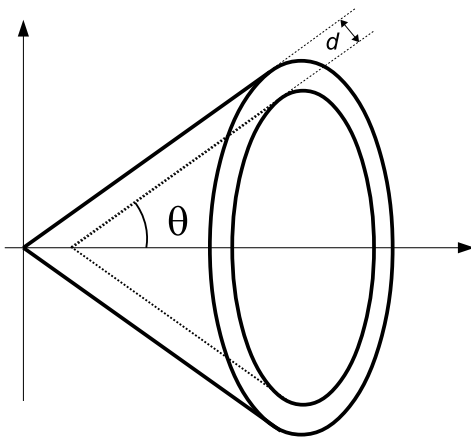


$\theta \neq \frac{\pi}{2}$ : Exner and Tater (2010): Dirichlet Laplacian  $A$  in a rotationally invariant infinite conical layer of constant width  $d$ .

Continuous spectrum =  $\left[ \frac{\pi^2}{d^2}, +\infty \right)$ , infinitely many eigenvalues below  $\frac{\pi^2}{d^2}$ .

# Previous works

Conical layer:



$\theta \neq \frac{\pi}{2}$ : Exner and Tater (2010): Dirichlet Laplacian  $A$  in a rotationally invariant infinite conical layer of constant width  $d$ .

Continuous spectrum =  $\left[ \frac{\pi^2}{d^2}, +\infty \right)$ , infinitely many eigenvalues below  $\frac{\pi^2}{d^2}$ .

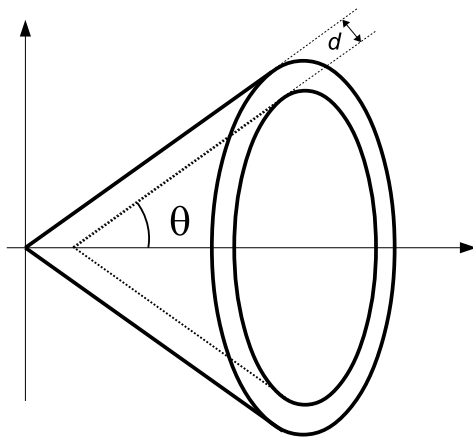
Dauge, Ourmières-Bonafos, Raymond (2015):

$$N(A, \frac{\pi^2}{d^2} - \varepsilon) \sim \frac{|\cot \theta|}{4\pi} |\ln \varepsilon|, \quad \varepsilon \rightarrow 0.$$



# Previous works

## Conical layer:



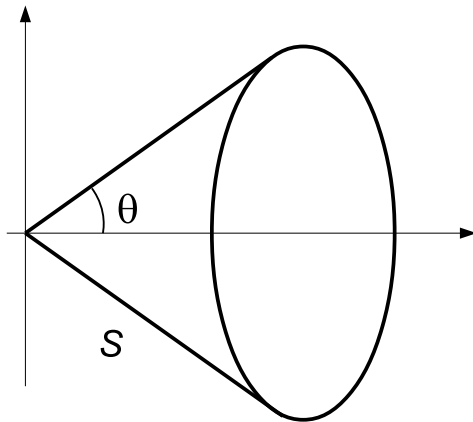
$\theta \neq \frac{\pi}{2}$ : Exner and Tater (2010): Dirichlet Laplacian  $A$  in a rotationally invariant infinite conical layer of constant width  $d$ .

Continuous spectrum =  $\left[\frac{\pi^2}{d^2}, +\infty\right)$ , infinitely many eigenvalues below  $\frac{\pi^2}{d^2}$ .

Dauge, Ourmières-Bonafos, Raymond (2015):

$$N(A, \frac{\pi^2}{d^2} - \varepsilon) \sim \frac{|\cot \theta|}{4\pi} |\ln \varepsilon|, \quad \varepsilon \rightarrow 0.$$

## $\delta$ -potential supported by cone:



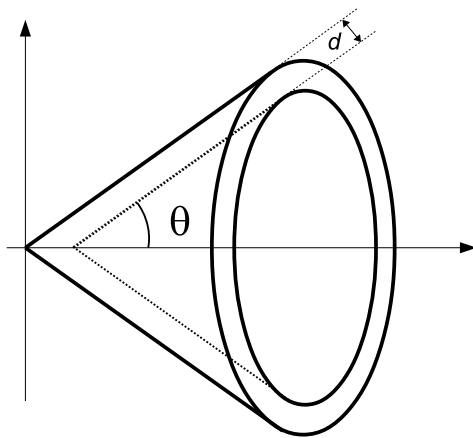
Let  $A$  be the Laplacian with the jump  $[\partial u / \partial n] = \alpha u$  at  $S$ ; the quadratic form is

$$H^1(\mathbb{R}^3) \ni u \mapsto \int_{\mathbb{R}^3} |\nabla u|^2 dx - \alpha \int_S |u|^2 ds, \quad \alpha > 0.$$

Continuous spectrum  $[-\alpha^2/4, +\infty)$ , infinitely many eigenvalues below  $(-\alpha^2/4)$ : Behrndt, Exner, Lotoreichik (2014).

# Previous works

## Conical layer:



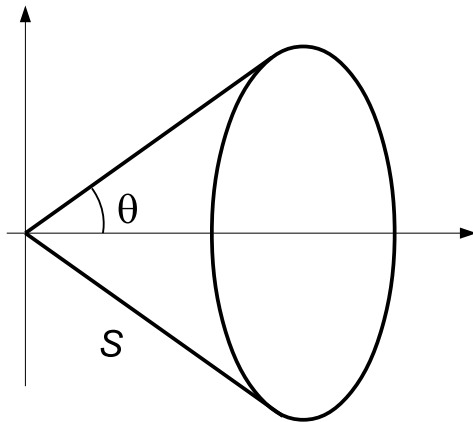
$\theta \neq \frac{\pi}{2}$ : Exner and Tater (2010): Dirichlet Laplacian  $A$  in a rotationally invariant infinite conical layer of constant width  $d$ .

Continuous spectrum =  $\left[\frac{\pi^2}{d^2}, +\infty\right)$ , infinitely many eigenvalues below  $\frac{\pi^2}{d^2}$ .

Dauge, Ourmières-Bonafos, Raymond (2015):

$$N(A, \frac{\pi^2}{d^2} - \varepsilon) \sim \frac{|\cot \theta|}{4\pi} |\ln \varepsilon|, \quad \varepsilon \rightarrow 0.$$

## $\delta$ -potential supported by cone:



Let  $A$  be the Laplacian with the jump  $[\partial u / \partial n] = \alpha u$  at  $S$ ; the quadratic form is

$$H^1(\mathbb{R}^3) \ni u \mapsto \int_{\mathbb{R}^3} |\nabla u|^2 dx - \alpha \int_S |u|^2 ds, \quad \alpha > 0.$$

Continuous spectrum  $[-\alpha^2/4, +\infty)$ , infinitely many eigenvalues below  $(-\alpha^2/4)$ : Behrndt, Exner, Lotoreichik (2014).

Lotoreichik, Ourmières-Bonafos (preprint 2015): the same asymptotics for  $N(A, -\alpha^2/4 - \varepsilon)$ .

# Geometry

**Further models? Arbitrarily shaped cones?**

# Geometry

## Further models? Arbitrarily shaped cones?

We say that a surface/domain  $X \subset \mathbb{R}^3$  is *conical* if

$$\lambda X = X \text{ for all } \lambda > 0.$$

Uniquely determined by the *spherical cross-section*  $X \cap \{|x| = 1\}$ .

# Geometry

## Further models? Arbitrarily shaped cones?

We say that a surface/domain  $X \subset \mathbb{R}^3$  is *conical* if

$$\lambda X = X \text{ for all } \lambda > 0.$$

Uniquely determined by the *spherical cross-section*  $X \cap \{|x| = 1\}$ .

If  $S$  is a conical surface, its  $\frac{d}{2}$ -neighborhood will be called *conical layer* of constant width  $d$ .

# Geometry

## Further models? Arbitrarily shaped cones?

We say that a surface/domain  $X \subset \mathbb{R}^3$  is *conical* if

$$\lambda X = X \text{ for all } \lambda > 0.$$

Uniquely determined by the *spherical cross-section*  $X \cap \{|x| = 1\}$ .

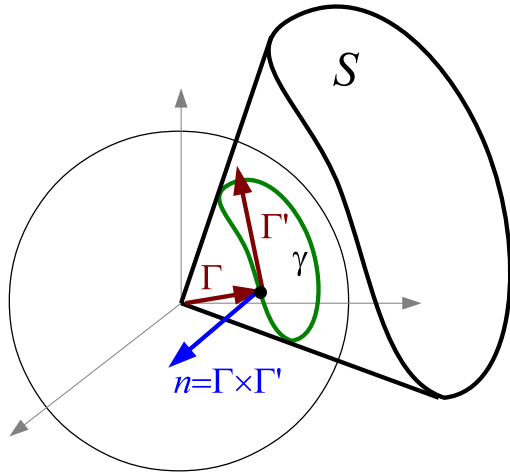
If  $S$  is a conical surface, its  $\frac{d}{2}$ -neighborhood will be called *conical layer* of constant width  $d$ .

## Three models:

- $\Omega$  is a conical layer, Dirichlet Laplacian in  $\Omega$ ,
- $S$  is a conical surface, Schrödinger operator with a  $\delta$ -potential on  $S$ ,
- $\Omega$  is a conical domain, the Laplacian in  $\Omega$  with Robin boundary condition

$$\frac{\partial u}{\partial n} = \alpha u.$$

# Curvature and layers

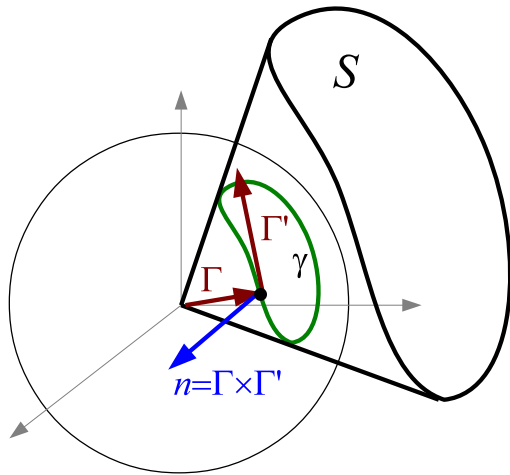


$S$  conical surface, **the cross-section**  $\gamma \subset \mathbb{S}^2$  is a **simple  $C^4$  loop** with an arc-length parametrization

$$\Gamma : \mathbb{T} \rightarrow \mathbb{R}^3, \quad \mathbb{T} := \mathbb{R}/(\ell\mathbb{Z}).$$

The tangent vector  $\Gamma'$ , the normal  $n = \Gamma \times \Gamma'$ .

# Curvature and layers



$S$  conical surface, **the cross-section**  $\gamma \subset \mathbb{S}^2$  is a **simple  $C^4$  loop** with an arc-length parametrization

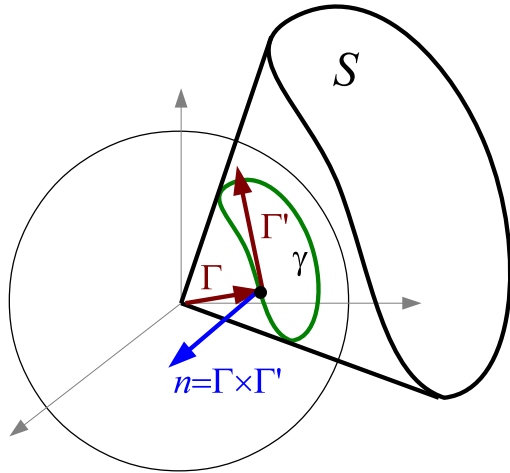
$$\Gamma : \mathbb{T} \rightarrow \mathbb{R}^3, \quad \mathbb{T} := \mathbb{R}/(\ell\mathbb{Z}).$$

The tangent vector  $\Gamma'$ , the normal  $n = \Gamma \times \Gamma'$ .

Geodesic curvature  $\kappa := [\Gamma'', \Gamma', \Gamma]$ .



# Curvature and layers



$S$  conical surface, **the cross-section**  $\gamma \subset \mathbb{S}^2$  is a **simple  $C^4$  loop** with an arc-length parametrization

$$\Gamma : \mathbb{T} \rightarrow \mathbb{R}^3, \quad \mathbb{T} := \mathbb{R}/(\ell\mathbb{Z}).$$

The tangent vector  $\Gamma'$ , the normal  $n = \Gamma \times \Gamma'$ .

Geodesic curvature  $\kappa := [\Gamma'', \Gamma', \Gamma]$ .

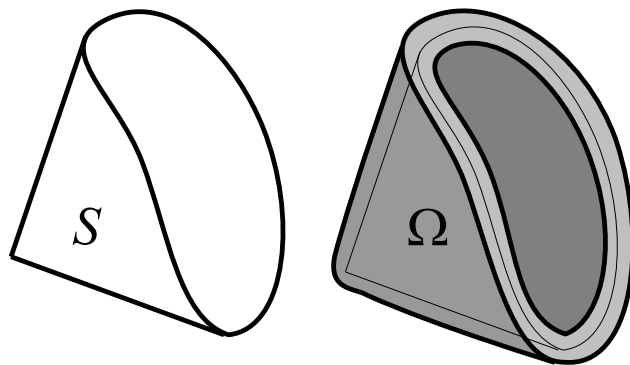
**Theorem.** Let  $\Omega$  be a conical layer of width  $d$  around  $S$  and  $A$  be the Dirichlet Laplacian in  $\Omega$ , then

$$N\left(A, \frac{\pi^2}{d^2} - \varepsilon\right) \sim C|\ln \varepsilon|, \quad \varepsilon \rightarrow 0, \quad C = \sum_{j=1}^{\infty} \sqrt{(-\mu_j)_+},$$

where  $\mu_j$  are the eigenvalues of

$$-\frac{d^2}{ds^2} - \frac{\kappa^2}{4} \text{ on } \mathbb{T}.$$

If  $\gamma$  is not a big circle ( $\sim S$  not a plane), then  $C > 0$  and there are infinitely many eigenvalues.



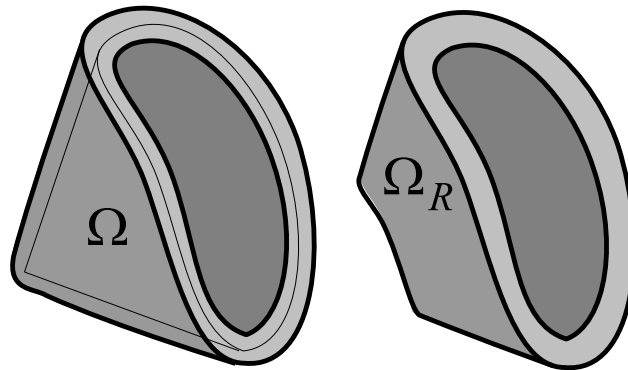
# Proof: Adapted coordinates

The surface  $S$  is parametrized by  $\mathbb{R}_+ \times \mathbb{T} \ni (r, s) \rightarrow r\Gamma(s)$ .

# Proof: Adapted coordinates

The surface  $S$  is parametrized by  $\mathbb{R}_+ \times \mathbb{T} \ni (r, s) \rightarrow r\Gamma(s)$ . The associated conical layer  $\Omega$  coincides, outside  $B_R(0)$ , with

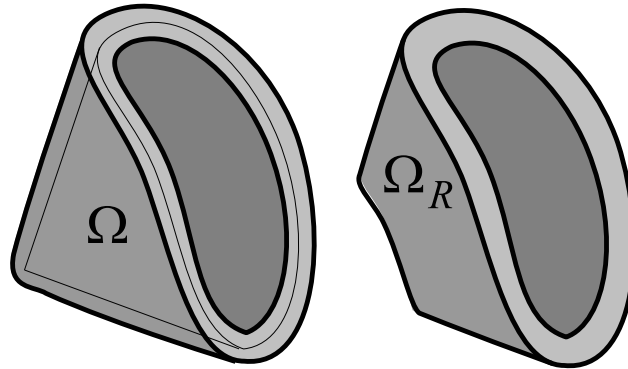
$$\Omega_R := X\left(\left(R, +\infty\right) \times \mathbb{T} \times \left(-\frac{d}{2}, \frac{d}{2}\right)\right), \quad X(r, s, t) = r\Gamma(s) + t\Gamma(s) \times \Gamma'(s).$$



# Proof: Adapted coordinates

The surface  $S$  is parametrized by  $\mathbb{R}_+ \times \mathbb{T} \ni (r, s) \rightarrow r\Gamma(s)$ . The associated conical layer  $\Omega$  coincides, outside  $B_R(0)$ , with

$$\Omega_R := X \left( (R, +\infty) \times \mathbb{T} \times \left( -\frac{d}{2}, \frac{d}{2} \right) \right), \quad X(r, s, t) = r\Gamma(s) + t\Gamma(s) \times \Gamma'(s).$$



The Dirichlet Laplacian  $A_\Omega$  in  $\Omega$  satisfies

$$\text{continuous spectrum} = \left[ \frac{\pi^2}{d^2}, +\infty \right),$$

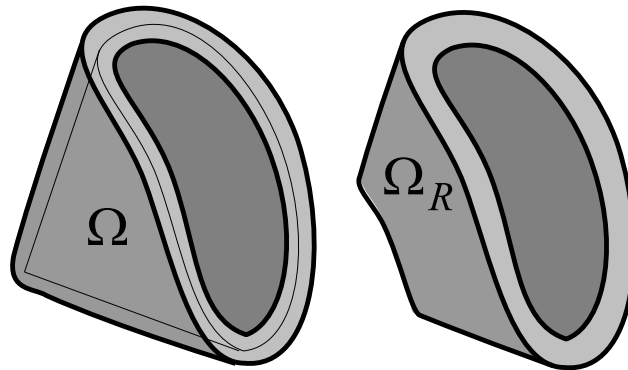
$$N \left( A_{R,D}, \frac{\pi^2}{d^2} - \varepsilon \right) - C \leq N \left( A_\Omega, \frac{\pi^2}{d^2} - \varepsilon \right) \leq N \left( A_{R,N}, \frac{\pi^2}{d^2} - \varepsilon \right) + C,$$

where  $A_{R,D/N}$  are Laplacians in  $\Omega_R$ , Dirichlet on  $\partial\Omega \cap \partial\Omega_R$ , Dirichlet/Neumann at  $|x| = R$ .

# Proof: Adapted coordinates

The surface  $S$  is parametrized by  $\mathbb{R}_+ \times \mathbb{T} \ni (r, s) \rightarrow r\Gamma(s)$ . The associated conical layer  $\Omega$  coincides, outside  $B_R(0)$ , with

$$\Omega_R := X \left( (R, +\infty) \times \mathbb{T} \times \left( -\frac{d}{2}, \frac{d}{2} \right) \right), \quad X(r, s, t) = r\Gamma(s) + t\Gamma(s) \times \Gamma'(s).$$



The Dirichlet Laplacian  $A_\Omega$  in  $\Omega$  satisfies

$$\text{continuous spectrum} = \left[ \frac{\pi^2}{d^2}, +\infty \right),$$

$$N\left(A_{R,D}, \frac{\pi^2}{d^2} - \varepsilon\right) - C \leq N\left(A_\Omega, \frac{\pi^2}{d^2} - \varepsilon\right) \leq N\left(A_{R,N}, \frac{\pi^2}{d^2} - \varepsilon\right) + C,$$

where  $A_{R,D/N}$  are Laplacians in  $\Omega_R$ , Dirichlet on  $\partial\Omega \cap \partial\Omega_R$ , Dirichlet/Neumann at  $|x| = R$ .

The asymptotics of  $N\left(A_{R,D/N}, \frac{\pi^2}{d^2} - \varepsilon\right)$  for  $\varepsilon \rightarrow 0$ ?

# Proof: 1D reduction

One passes to  $L^2(\Pi, dr ds dt)$ ,  $\Pi := (R, +\infty) \times \mathbb{T} \times \left(-\frac{d}{2}, \frac{d}{2}\right)$ .

After some transformations, one minorates/majorates the quadratic forms for  $A := A_{R,N/D}$  in  $(r, s, t)$  by

$$\int_{\Pi} \left( v_r^2 + \frac{1}{r^2} \left( v_s^2 - \frac{\kappa^2 + 1}{4} v^2 \right) + v_t^2 + \frac{c}{r^3} v^2 \right) dr ds dt + c \int_{r=R} v^2 ds dt, \quad v(r, \cdot) = 0 \text{ for } D.$$

# Proof: 1D reduction

One passes to  $L^2(\Pi, dr ds dt)$ ,  $\Pi := (R, +\infty) \times \mathbb{T} \times \left(-\frac{d}{2}, \frac{d}{2}\right)$ .

After some transformations, one minorates/majorates the quadratic forms for  $A := A_{R,N/D}$  in  $(r, s, t)$  by

$$\int_{\Pi} \left( v_r^2 + \frac{1}{r^2} \left( v_s^2 - \frac{\kappa^2 + 1}{4} v^2 \right) + v_t^2 + \frac{c}{r^3} v^2 \right) dr ds dt + c \int_{r=R} v^2 ds dt, \quad v(r, \cdot) = 0 \text{ for } D.$$

One represents

$$v(r, s, t) = \sum_{j,k} f_{j,k}(r) g_j(s) h_k(s),$$

where

# Proof: 1D reduction

One passes to  $L^2(\Pi, dr ds dt)$ ,  $\Pi := (R, +\infty) \times \mathbb{T} \times \left(-\frac{d}{2}, \frac{d}{2}\right)$ .

After some transformations, one minorates/majorates the quadratic forms for  $A := A_{R,N/D}$  in  $(r, s, t)$  by

$$\int_{\Pi} \left( v_r^2 + \frac{1}{r^2} \left( v_s^2 - \frac{\kappa^2 + 1}{4} v^2 \right) + v_t^2 + \frac{c}{r^3} v^2 \right) dr ds dt + c \int_{r=R} v^2 ds dt, \quad v(r, \cdot) = 0 \text{ for } D.$$

One represents

$$v(r, s, t) = \sum_{j,k} f_{j,k}(r) g_j(s) h_k(s),$$

where

- $h_k$  are the normalized eigenfunctions of  $h \mapsto -h''$  on  $\left(-\frac{d}{2}, \frac{d}{2}\right)$  with the Dirichlet boundary conditions. The eigenvalues are  $\frac{\pi^2 k^2}{d^2}$ ,  $k = 1, 2, \dots$ ,



# Proof: 1D reduction

One passes to  $L^2(\Pi, dr ds dt)$ ,  $\Pi := (R, +\infty) \times \mathbb{T} \times \left(-\frac{d}{2}, \frac{d}{2}\right)$ .

After some transformations, one minorates/majorates the quadratic forms for  $A := A_{R,N/D}$  in  $(r, s, t)$  by

$$\int_{\Pi} \left( v_r^2 + \frac{1}{r^2} \left( v_s^2 - \frac{\kappa^2 + 1}{4} v^2 \right) + v_t^2 + \frac{c}{r^3} v^2 \right) dr ds dt + c \int_{r=R} v^2 ds dt, \quad v(r, \cdot) = 0 \text{ for } D.$$

One represents

$$v(r, s, t) = \sum_{j,k} f_{j,k}(r) g_j(s) h_k(s),$$

where

- $h_k$  are the normalized eigenfunctions of  $h \mapsto -h''$  on  $\left(-\frac{d}{2}, \frac{d}{2}\right)$  with the Dirichlet boundary conditions. The eigenvalues are  $\frac{\pi^2 k^2}{d^2}$ ,  $k = 1, 2, \dots$ ,
- $g_j$  are the orthonormalized eigenfunctions of  $g \mapsto -g'' - \frac{\kappa^2}{4} g$  on  $\mathbb{T}$ . The eigenvalues are  $\mu_j$ .

# Proof: 1D reduction

One passes to  $L^2(\Pi, dr ds dt)$ ,  $\Pi := (R, +\infty) \times \mathbb{T} \times \left(-\frac{d}{2}, \frac{d}{2}\right)$ .

After some transformations, one minorates/majorates the quadratic forms for  $A := A_{R,N/D}$  in  $(r, s, t)$  by

$$\int_{\Pi} \left( v_r^2 + \frac{1}{r^2} \left( v_s^2 - \frac{\kappa^2 + 1}{4} v^2 \right) + v_t^2 + \frac{c}{r^3} v^2 \right) dr ds dt + c \int_{r=R} v^2 ds dt, \quad v(r, \cdot) = 0 \text{ for } D.$$

One represents

$$v(r, s, t) = \sum_{j,k} f_{j,k}(r) g_j(s) h_k(s),$$

where

- $h_k$  are the normalized eigenfunctions of  $h \mapsto -h''$  on  $\left(-\frac{d}{2}, \frac{d}{2}\right)$  with the Dirichlet boundary conditions. The eigenvalues are  $\frac{\pi^2 k^2}{d^2}$ ,  $k = 1, 2, \dots$ ,
- $g_j$  are the orthonormalized eigenfunctions of  $g \mapsto -g'' - \frac{\kappa^2}{4} g$  on  $\mathbb{T}$ . The eigenvalues are  $\mu_j$ .

Then  $A \simeq \bigoplus_{j,k} A_{j,k}$  with  $A_{j,k} = -\frac{d^2}{dr^2} + \frac{\mu_j - 1/4}{r^2} + \frac{c}{r^3} + \left(\frac{\pi k}{d}\right)^2$  in  $L^2(R, \infty)$  (+ b.c. at  $R$ ).

# Proof: 1D reduction

One passes to  $L^2(\Pi, dr ds dt)$ ,  $\Pi := (R, +\infty) \times \mathbb{T} \times \left(-\frac{d}{2}, \frac{d}{2}\right)$ .

After some transformations, one minorates/majorates the quadratic forms for  $A := A_{R,N/D}$  in  $(r, s, t)$  by

$$\int_{\Pi} \left( v_r^2 + \frac{1}{r^2} \left( v_s^2 - \frac{\kappa^2 + 1}{4} v^2 \right) + v_t^2 + \frac{c}{r^3} v^2 \right) dr ds dt + c \int_{r=R} v^2 ds dt, \quad v(r, \cdot) = 0 \text{ for } D.$$

One represents

$$v(r, s, t) = \sum_{j,k} f_{j,k}(r) g_j(s) h_k(s),$$

where

- $h_k$  are the normalized eigenfunctions of  $h \mapsto -h''$  on  $\left(-\frac{d}{2}, \frac{d}{2}\right)$  with the Dirichlet boundary conditions. The eigenvalues are  $\frac{\pi^2 k^2}{d^2}$ ,  $k = 1, 2, \dots$ ,
- $g_j$  are the orthonormalized eigenfunctions of  $g \mapsto -g'' - \frac{\kappa^2}{4} g$  on  $\mathbb{T}$ . The eigenvalues are  $\mu_j$ .

Then  $A \simeq \bigoplus_{j,k} A_{j,k}$  with  $A_{j,k} = -\frac{d^2}{dr^2} + \frac{\mu_j - 1/4}{r^2} + \frac{c}{r^3} + \left(\frac{\pi k}{d}\right)^2$  in  $L^2(R, \infty)$  (+ b.c. at  $R$ ).

Only  $k = 1$  and  $\mu_j < 0$  contribute, and one can use the known 1D result.

# $\delta$ -interaction

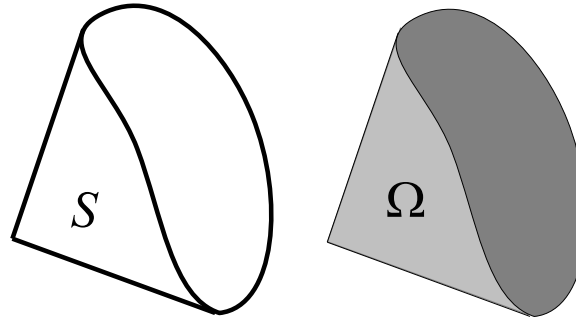
Under the same assumptions on  $S$ , literally the same result is expected for the Schrödinger operators with  $\delta$ -potentials: the associated quadratic form is

$$H^1(\mathbb{R}^3) \ni u \mapsto \int_{\mathbb{R}^3} |\nabla u|^2 dx - \alpha \int_S |u|^2 ds.$$

Some technical details are still to be worked out: the analysis on a conical layer of a non-constant width around  $S$  (in preparation, with T. Ourmières-Bonafos)

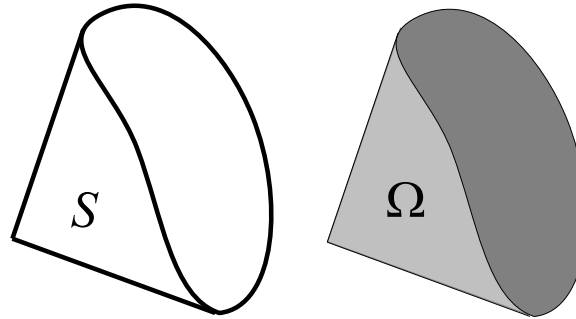
# Robin Laplacian

Let  $\Omega$  be bounded by a conical surface  $S$  as above.



# Robin Laplacian

Let  $\Omega$  be bounded by a conical surface  $S$  as above.



For  $\alpha > 0$ , consider  $Q_\alpha^\Omega$  Laplacian in  $L^2(\Omega)$  with Robin boundary conditions:

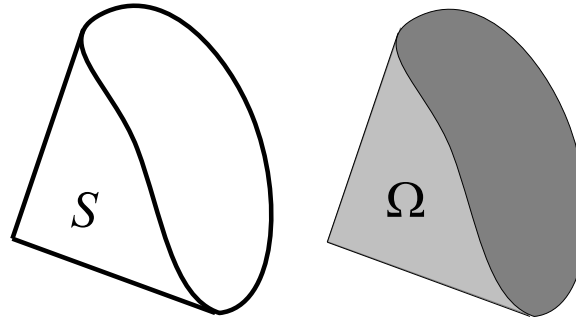
$$u \mapsto -\Delta u, \quad \frac{\partial u}{\partial n} = \alpha u \text{ at } S.$$

The quadratic form is

$$H^1(\Omega) \ni u \mapsto \int_{\Omega} |\nabla u|^2 dx - \alpha \int_S |u|^2 ds.$$

# Robin Laplacian

Let  $\Omega$  be bounded by a conical surface  $S$  as above.



For  $\alpha > 0$ , consider  $Q_\alpha^\Omega$  Laplacian in  $L^2(\Omega)$  with Robin boundary conditions:

$$u \mapsto -\Delta u, \quad \frac{\partial u}{\partial n} = \alpha u \text{ at } S.$$

The quadratic form is

$$H^1(\Omega) \ni u \mapsto \int_{\Omega} |\nabla u|^2 dx - \alpha \int_S |u|^2 ds.$$

Difference to the previous cases:

**Theorem (KP'2015).** The continuous spectrum is  $[-\alpha^2, \infty)$  and:

- If  $\Omega$  is the exterior of a convex set, then the discrete spectrum is empty.
- Otherwise, infinite many eigenvalues below  $-\alpha^2$ .

(Attention: smoothness of the boundary is important!)

# Proof for convex $\Omega^c$

Continuous spectrum  $[-\alpha^2, +\infty)$ : standard (cutting out a compact part + approximate eigenfunctions): Bruneau, Popoff (2016).



# Proof for convex $\Omega^c$

Continuous spectrum  $[-\alpha^2, +\infty)$ : standard (cutting out a compact part + approximate eigenfunctions): Bruneau, Popoff (2016).

If  $\Omega^c$  is convex, then  $\Phi(S \times \mathbb{R}_+) \subset \Omega$ , with  $\Phi(s, t) = s - tn(s)$  injective.

# Proof for convex $\Omega^c$

Continuous spectrum  $[-\alpha^2, +\infty)$ : standard (cutting out a compact part + approximate eigenfunctions): Bruneau, Popoff (2016).

If  $\Omega^c$  is convex, then  $\Phi(S \times \mathbb{R}_+) \subset \Omega$ , with  $\Phi(s, t) = s - tn(s)$  injective. One has, with  $J(s, t) = \prod (1 - tk_j(s))$ ,  $k_j$  the principal curvatures of  $S$ ,

$$\begin{aligned} & \int_{\Omega} |\nabla u|^2 ds - \alpha \int_S |u|^2 ds + \alpha^2 \int_{\Omega} |u|^2 dx \\ & \geq \int_{\Phi(S \times \mathbb{R}_+)} |\nabla u|^2 ds - \alpha \int_S |u|^2 ds + \alpha^2 \int_{\Phi(S \times \mathbb{R}_+)} |u|^2 dx \\ & = \int_{S \times \mathbb{R}_+} \left| (\nabla u)(\Phi(s, t)) \right|^2 J(s, t) ds dt - \alpha \int_S |u|^2 ds + \alpha^2 \int_{S \times \mathbb{R}_+} \left| u(P(s, t)) \right|^2 J(s, t) ds dt =: I. \end{aligned}$$

# Proof for convex $\Omega^c$

Continuous spectrum  $[-\alpha^2, +\infty)$ : standard (cutting out a compact part + approximate eigenfunctions): Bruneau, Popoff (2016).

If  $\Omega^c$  is convex, then  $\Phi(S \times \mathbb{R}_+) \subset \Omega$ , with  $\Phi(s, t) = s - tn(s)$  injective. One has, with  $J(s, t) = \prod (1 - tk_j(s))$ ,  $k_j$  the principal curvatures of  $S$ ,

$$\begin{aligned} & \int_{\Omega} |\nabla u|^2 ds - \alpha \int_S |u|^2 ds + \alpha^2 \int_{\Omega} |u|^2 dx \\ & \geq \int_{\Phi(S \times \mathbb{R}_+)} |\nabla u|^2 ds - \alpha \int_S |u|^2 ds + \alpha^2 \int_{\Phi(S \times \mathbb{R}_+)} |u|^2 dx \\ & = \int_{S \times \mathbb{R}_+} \left| (\nabla u)(\Phi(s, t)) \right|^2 J(s, t) ds dt - \alpha \int_S |u|^2 ds + \alpha^2 \int_{S \times \mathbb{R}_+} \left| u(P(s, t)) \right|^2 J(s, t) ds dt =: I. \end{aligned}$$

Due to  $k_j \leq 0$  we have  $J \geq 1$ , and

$$\left| (\nabla u)(\Phi(s, t)) \right|^2 \geq \left| (n \cdot \nabla u)(\Phi(s, t)) \right|^2 = \left| \frac{\partial}{\partial t} u(\Phi(s, t)) \right|^2,$$

# Proof for convex $\Omega^c$

Continuous spectrum  $[-\alpha^2, +\infty)$ : standard (cutting out a compact part + approximate eigenfunctions): Bruneau, Popoff (2016).

If  $\Omega^c$  is convex, then  $\Phi(S \times \mathbb{R}_+) \subset \Omega$ , with  $\Phi(s, t) = s - tn(s)$  injective. One has, with  $J(s, t) = \prod (1 - tk_j(s))$ ,  $k_j$  the principal curvatures of  $S$ ,

$$\begin{aligned} & \int_{\Omega} |\nabla u|^2 ds - \alpha \int_S |u|^2 ds + \alpha^2 \int_{\Omega} |u|^2 dx \\ & \geq \int_{\Phi(S \times \mathbb{R}_+)} |\nabla u|^2 ds - \alpha \int_S |u|^2 ds + \alpha^2 \int_{\Phi(S \times \mathbb{R}_+)} |u|^2 dx \\ & = \int_{S \times \mathbb{R}_+} \left| (\nabla u)(\Phi(s, t)) \right|^2 J(s, t) ds dt - \alpha \int_S |u|^2 ds + \alpha^2 \int_{S \times \mathbb{R}_+} \left| u(P(s, t)) \right|^2 J(s, t) ds dt =: I. \end{aligned}$$

Due to  $k_j \leq 0$  we have  $J \geq 1$ , and

$$\left| (\nabla u)(\Phi(s, t)) \right|^2 \geq \left| (n \cdot \nabla u)(\Phi(s, t)) \right|^2 = \left| \frac{\partial}{\partial t} u(\Phi(s, t)) \right|^2,$$

and

$$I \geq \int_S \left( \int_{\mathbb{R}_+} \left| \frac{\partial}{\partial t} u(\Phi(s, t)) \right|^2 dt - \alpha |u(s, 0)|^2 + \alpha^2 \int_{\mathbb{R}_+} |u(\Phi(s, t))|^2 dt \right) ds \geq 0,$$

i.e.  $Q_{\alpha}^{\Omega} > -\alpha^2$ .

# Counting the eigenvalues

**Theorem (Bruneau, KP, Popoff, 2016).**

$$N(Q_\alpha^\Omega, -\alpha^2 - \varepsilon) \sim \frac{\alpha^2}{8\pi\varepsilon} \int_{\mathbb{T}} (\kappa_+)^2 ds, \quad \varepsilon \rightarrow 0. \quad (\star)$$

(Note:  $\Omega^c$  convex iff  $\kappa \leq 0$ )

# Counting the eigenvalues

**Theorem (Bruneau, KP, Popoff, 2016).**

$$N(Q_\alpha^\Omega, -\alpha^2 - \varepsilon) \sim \frac{\alpha^2}{8\pi\varepsilon} \int_{\mathbb{T}} (\kappa_+)^2 ds, \quad \varepsilon \rightarrow 0. \quad (\star)$$

(Note:  $\Omega^c$  convex iff  $\kappa \leq 0$ )

Context (Pankrashkin, Popoff, 2014–2015): for smooth  $\Omega$ ,

$$E_j(Q_\alpha^\Omega) \sim -\alpha^2 + E_j(-\Delta_S - \alpha H) + O(1) \text{ as } \alpha \rightarrow +\infty,$$

$\Delta_S$  and  $H$  are Laplace-Beltrami operator and the mean curvature on  $S$  (then used by Kachmar, Keraval, Raymond for the eigenvalue counting).

# Counting the eigenvalues

**Theorem (Bruneau, KP, Popoff, 2016).**

$$N(Q_\alpha^\Omega, -\alpha^2 - \varepsilon) \sim \frac{\alpha^2}{8\pi\varepsilon} \int_{\mathbb{T}} (\kappa_+)^2 ds, \quad \varepsilon \rightarrow 0. \quad (\star)$$

(Note:  $\Omega^c$  convex iff  $\kappa \leq 0$ )

Context (Pankrashkin, Popoff, 2014–2015): for smooth  $\Omega$ ,

$$E_j(Q_\alpha^\Omega) \sim -\alpha^2 + E_j(-\Delta_S - \alpha H) + O(1) \text{ as } \alpha \rightarrow +\infty,$$

$\Delta_S$  and  $H$  are Laplace-Beltrami operator and the mean curvature on  $S$  (then used by Kachmar, Keraval, Raymond for the eigenvalue counting).

Equation  $(\star)$  is (formally) the Weyl asymptotics for  $N(-\Delta_S - \alpha H, -\varepsilon)$  with

$$H(r, s) = \frac{\kappa(s)}{r}.$$

# Counting the eigenvalues

**Theorem (Bruneau, KP, Popoff, 2016).**

$$N(Q_\alpha^\Omega, -\alpha^2 - \varepsilon) \sim \frac{\alpha^2}{8\pi\varepsilon} \int_{\mathbb{T}} (\kappa_+)^2 ds, \quad \varepsilon \rightarrow 0. \quad (\star)$$

(Note:  $\Omega^c$  convex iff  $\kappa \leq 0$ )

Context (Pankrashkin, Popoff, 2014–2015): for smooth  $\Omega$ ,

$$E_j(Q_\alpha^\Omega) \sim -\alpha^2 + E_j(-\Delta_S - \alpha H) + O(1) \text{ as } \alpha \rightarrow +\infty,$$

$\Delta_S$  and  $H$  are Laplace-Beltrami operator and the mean curvature on  $S$  (then used by Kachmar, Keraval, Raymond for the eigenvalue counting).

Equation  $(\star)$  is (formally) the Weyl asymptotics for  $N(-\Delta_S - \alpha H, -\varepsilon)$  with

$$H(r, s) = \frac{\kappa(s)}{r}.$$

**But in our case  $\alpha$  is fixed!**



# Proof: Polar coordinates

Let  $\Sigma := \Omega \cap \{|x| = 1\}$  be the cross section (spherical domain), then

$$\mathbb{R}_+ \times \Sigma \ni (r, \theta) \mapsto r\theta$$

is a parametrization of  $\Omega$ .

# Proof: Polar coordinates

Let  $\Sigma := \Omega \cap \{|x| = 1\}$  be the cross section (spherical domain), then

$$\mathbb{R}_+ \times \Sigma \ni (r, \theta) \mapsto r\theta$$

is a parametrization of  $\Omega$ . In  $L^2(\Omega) \simeq L^2(\mathbb{R}_+) \otimes L^2(\Sigma)$ , the quadratic form writes as

$$\int_{\mathbb{R}_+} \left( |u_r|^2 + \frac{1}{r^2} b_r(u(r, \cdot)) \right) dr, \quad b_r(v) = \int_{\Sigma} |\nabla_{\theta} v|^2 d\theta - \alpha r \int_{\gamma} |v|^2 d\ell,$$

# Proof: Polar coordinates

Let  $\Sigma := \Omega \cap \{|x| = 1\}$  be the cross section (spherical domain), then

$$\mathbb{R}_+ \times \Sigma \ni (r, \theta) \mapsto r\theta$$

is a parametrization of  $\Omega$ . In  $L^2(\Omega) \simeq L^2(\mathbb{R}_+) \otimes L^2(\Sigma)$ , the quadratic form writes as

$$\int_{\mathbb{R}_+} \left( |u_r|^2 + \frac{1}{r^2} b_r(u(r, \cdot)) \right) dr, \quad b_r(v) = \int_{\Sigma} |\nabla_{\theta} v|^2 d\theta - \alpha r \int_{\gamma} |v|^2 d\ell,$$

i.e.  $b_r$  corresponds to “ $Q_{\alpha r}^{\Sigma}$ ” for spherical domains  $\sim \left( -\frac{\partial^2}{\partial s^2} - \alpha r \kappa \right)$  on  $\gamma$  (Born-Oppenheimer).

# Proof: Polar coordinates

Let  $\Sigma := \Omega \cap \{|x| = 1\}$  be the cross section (spherical domain), then

$$\mathbb{R}_+ \times \Sigma \ni (r, \theta) \mapsto r\theta$$

is a parametrization of  $\Omega$ . In  $L^2(\Omega) \simeq L^2(\mathbb{R}_+) \otimes L^2(\Sigma)$ , the quadratic form writes as

$$\int_{\mathbb{R}_+} \left( |u_r|^2 + \frac{1}{r^2} b_r(u(r, \cdot)) \right) dr, \quad b_r(v) = \int_{\Sigma} |\nabla_{\theta} v|^2 d\theta - \alpha r \int_{\gamma} |v|^2 d\ell,$$

i.e.  $b_r$  corresponds to “ $Q_{\alpha r}^{\Sigma}$ ” for spherical domains  $\sim \left( -\frac{\partial^2}{\partial s^2} - \alpha r \kappa \right)$  on  $\gamma$  (Born-Oppenheimer).

**Theorem.** For some  $C > 0$ , there holds

$$N(L_1, -\varepsilon) - C \leq N(Q_{\alpha}^{\Omega}, -\alpha^2 - \varepsilon) \leq N(L_2, -\varepsilon) + C,$$

with  $L_j$  acting in  $L^2(\mathbb{R}_+ \times \mathbb{T})$  with the quadratic forms

$$u \mapsto \int_{\mathbb{R}_+} \int_{\mathbb{T}} \left( (1 + a_j(r)) |u_r|^2 + \frac{1 + b_j(r)}{r^2} |u_s|^2 - \frac{\alpha \kappa(s) + c_j(r)}{r} |u|^2 \right) ds dr,$$

$a_j, b_j, c_j$  vanishing at  $\infty$ , and Dirichlet/Neumann b.c. at  $r = 0$ .

# Proof: Polar coordinates

Let  $\Sigma := \Omega \cap \{|x| = 1\}$  be the cross section (spherical domain), then

$$\mathbb{R}_+ \times \Sigma \ni (r, \theta) \mapsto r\theta$$

is a parametrization of  $\Omega$ . In  $L^2(\Omega) \simeq L^2(\mathbb{R}_+) \otimes L^2(\Sigma)$ , the quadratic form writes as

$$\int_{\mathbb{R}_+} \left( |u_r|^2 + \frac{1}{r^2} b_r(u(r, \cdot)) \right) dr, \quad b_r(v) = \int_{\Sigma} |\nabla_{\theta} v|^2 d\theta - \alpha r \int_{\gamma} |v|^2 dl,$$

i.e.  $b_r$  corresponds to “ $Q_{\alpha r}^{\Sigma}$ ” for spherical domains  $\sim \left( -\frac{\partial^2}{\partial s^2} - \alpha r \kappa \right)$  on  $\gamma$  (Born-Oppenheimer).

**Theorem.** For some  $C > 0$ , there holds

$$N(L_1, -\varepsilon) - C \leq N(Q_{\alpha}^{\Omega}, -\alpha^2 - \varepsilon) \leq N(L_2, -\varepsilon) + C,$$

with  $L_j$  acting in  $L^2(\mathbb{R}_+ \times \mathbb{T})$  with the quadratic forms

$$u \mapsto \int_{\mathbb{R}_+} \int_{\mathbb{T}} \left( (1 + a_j(r)) |u_r|^2 + \frac{1 + b_j(r)}{r^2} |u_s|^2 - \frac{\alpha \kappa(s) + c_j(r)}{r} |u|^2 \right) ds dr,$$

$a_j, b_j, c_j$  vanishing at  $\infty$ , and Dirichlet/Neumann b.c. at  $r = 0$ .

Weyl asymptotics for  $N(L_j, -\varepsilon)$ : manual proof.