

Total absorption in waveguides by admittance conjugation

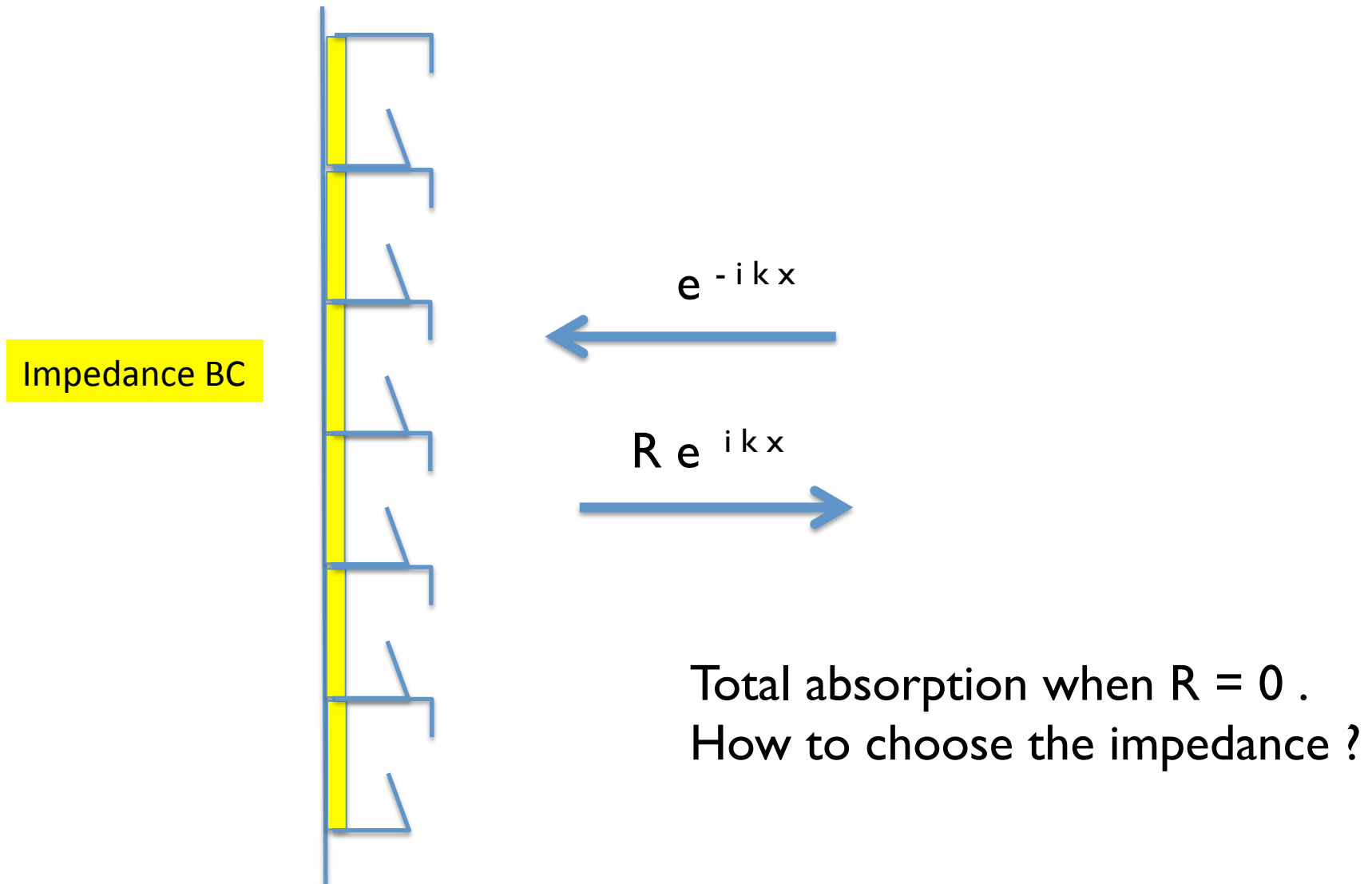
V. Pagneux

Lab. Acous. Univ. Maine – CNRS

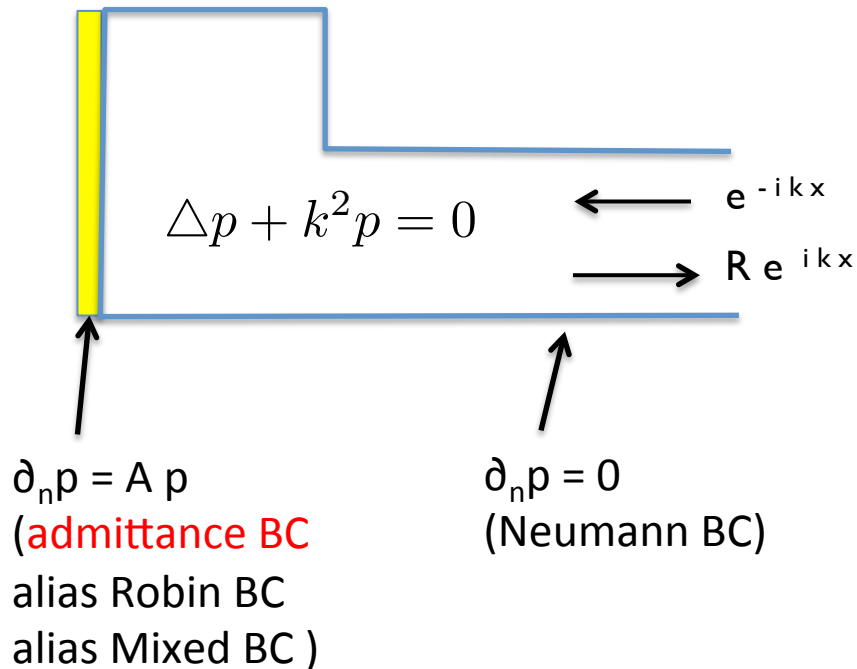
Le Mans - France

CONFERENCE ON WAVEGUIDES, Porquerolles
17th to 19th may 2016

Wall with total absorption



Will consider waveguide :
Reflection problem (acoustics)

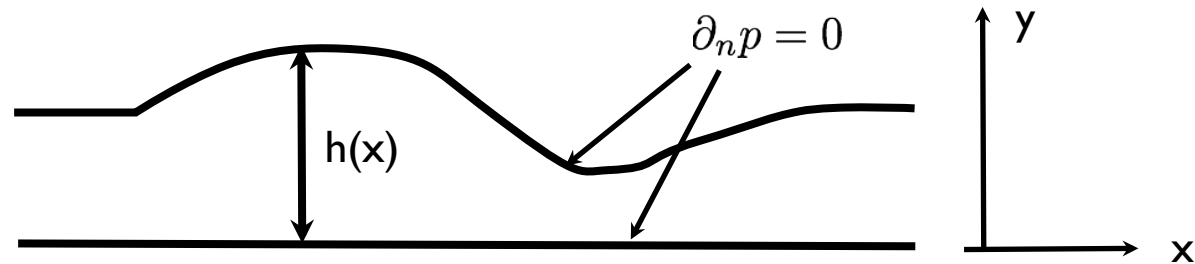


Absorption by admittance at the wall:
 $\text{Im}(A) < 0$ (convention $e^{-i\omega t}$)

- Looking for A values such that $R=0$ (total absorption)
- Difficult at low frequency : for small k , A has to be small ($|A| \ll k$ for thin coating)
- Surprising connection with water waves

coupled mode equations

2D acoustics



evolution equation

$$\partial_x^2 p + \partial_y^2 p + k^2 p = 0 \quad \Rightarrow \quad \partial_x \begin{pmatrix} p \\ \partial_x p \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -k^2 - \partial_y^2 & 0 \end{pmatrix} \begin{pmatrix} p \\ \partial_x p \end{pmatrix}$$

modal expansion

$$p(x, y) = \sum_{n \geq 0} a_n(x) g_n(y; x)$$

$$\partial_x p(x, y) = \sum_{n \geq 0} b_n(x) g_n(y; x)$$

local modes (Neumann)

$$g_n(y; x) = \sqrt{\frac{2 - \delta_{n0}}{h}} \cos(n\pi y/h)$$

coupled mode equations:

$$a' = -F a + b$$

$$b' = -K^2 a + F^T b$$

Unstable
evolution
eqs.

Riccati equation

admittance matrix
(~ DtN operator)

$$\underset{\partial_x p}{\uparrow} b = Y \underset{p}{\uparrow} a$$

from coupled mode eqs.

$$a' = -Fa + b$$

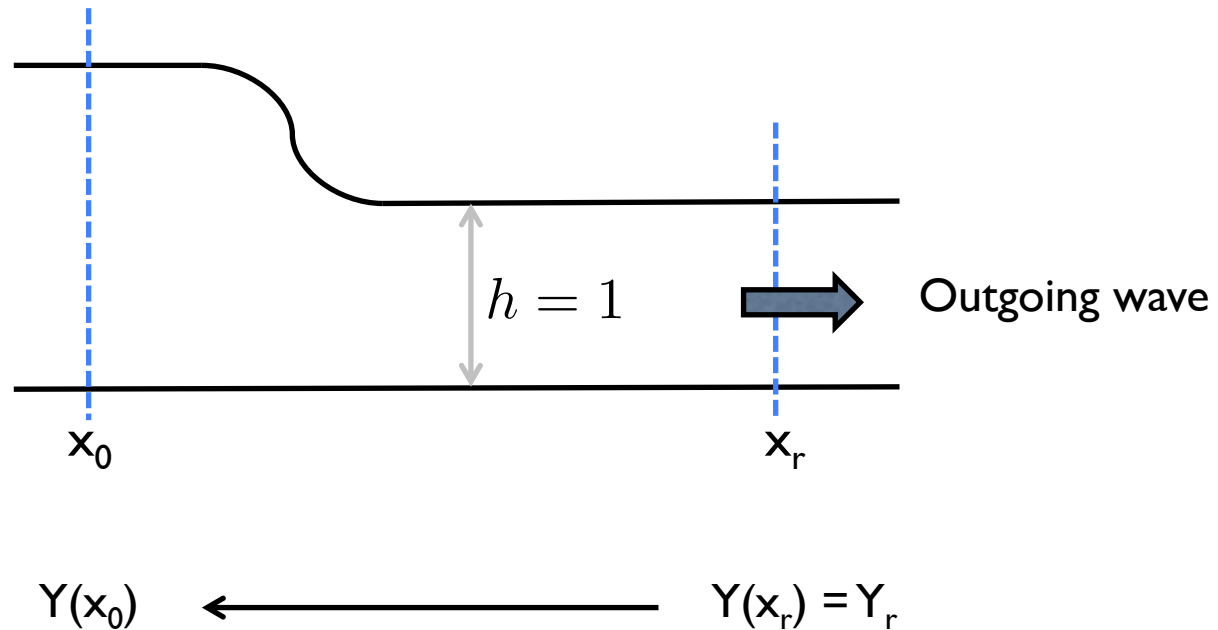
$$b' = -K^2 a + F^T b$$

$$Y' = -K^2 - Y^2 + YF + F^T Y$$

Riccati equation

- initial condition : radiation condition (e.g. outgoing wave)
- numerically stable
- the same for the impedance matrix Z
- gives the scattering matrix (reflection & transmission)

Admittance computation



$Y(x_0)$ known from Y_r by the Riccati equation

$$Y' = -K^2 - Y^2 + YF + F^T Y$$

Expression of the Y_r for outgoing waves

$$p = \sum_{n=0}^{\infty} a_n(x) g_n(y)$$

$$\partial_x p = \sum_{n=0}^{\infty} b_n(x) g_n(y)$$

$$a_n(x) = c_n e^{ik_n x} \longrightarrow b_n = ik_n a_n$$

Thus Y_r diagonal :

$$Y_r = \begin{pmatrix} ik_0 & & & \\ & ik_1 & & \\ & & ik_2 & \\ & & & \dots \end{pmatrix}$$

$$k_n = \sqrt{k^2 - n^2 \pi^2}$$

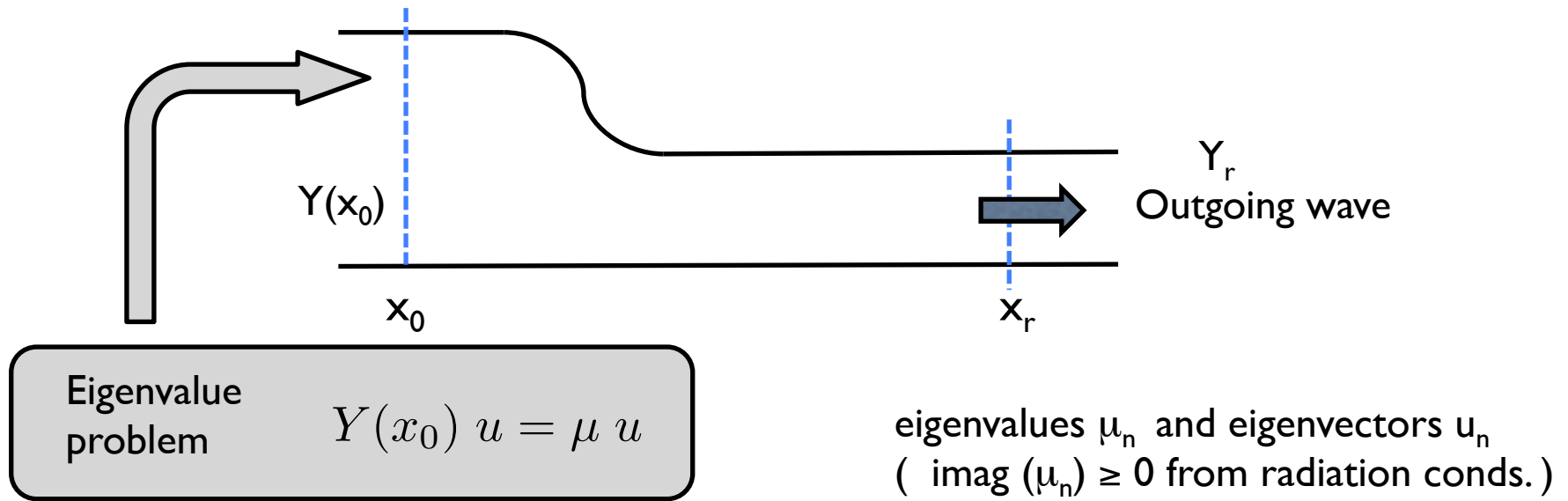
convention $e^{-i\omega t}$

k_n real (>0) : **propagating**

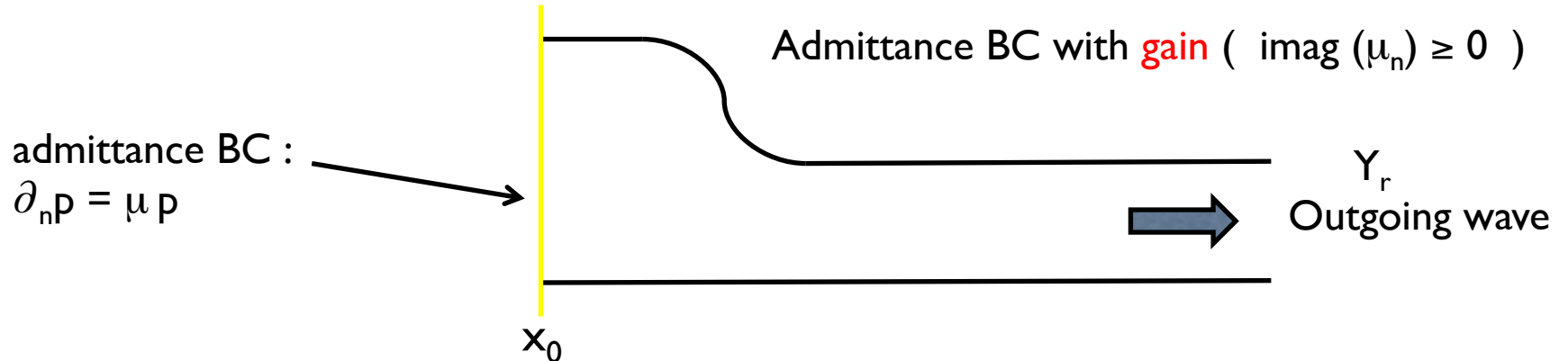
k_n pure imaginary ($\text{Imag} > 0$) : **evanescent**

$k_0 = k$: mode $n=0$ always propagating

Eigenvalues of $Y(x_0)$

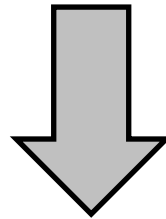
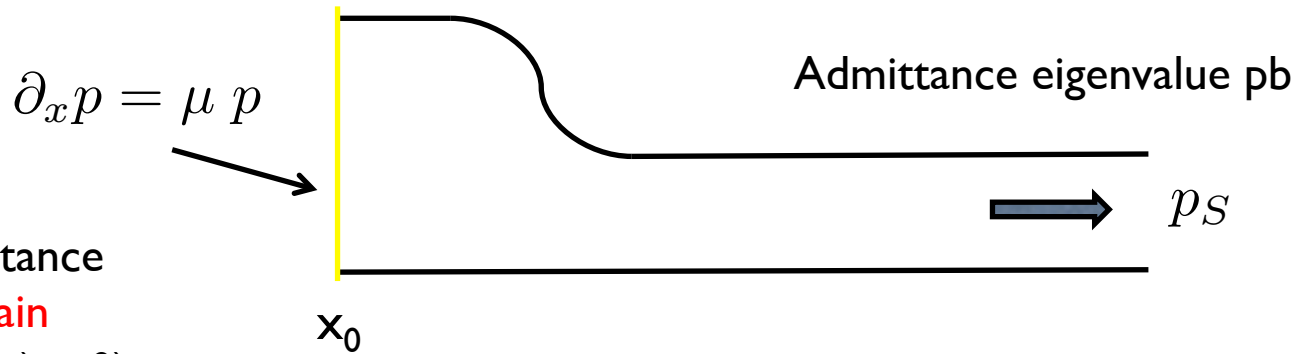


Solution associated to μ_n

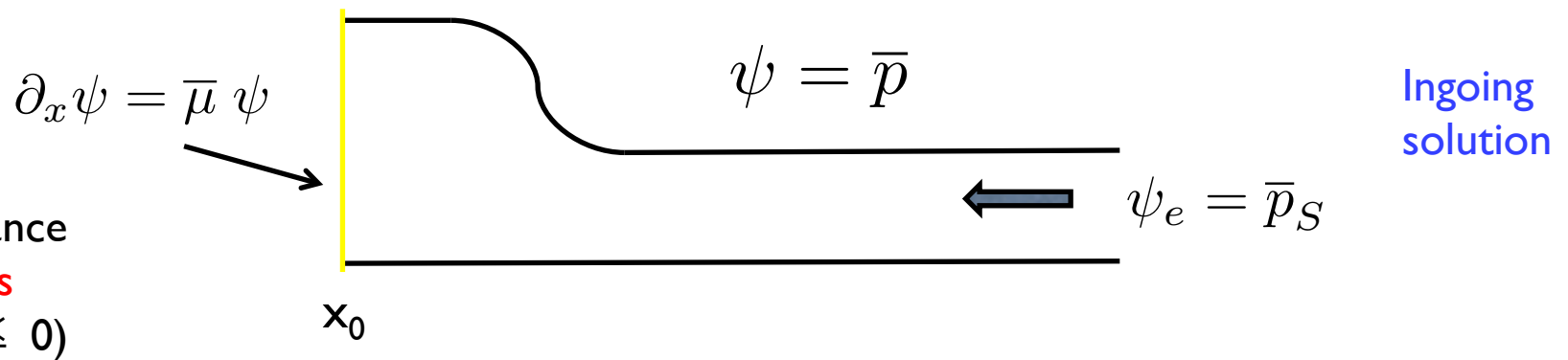


Remark: if $\mu=0$ **trapped mode** is found

Admittance conjugation



Phase conjugation



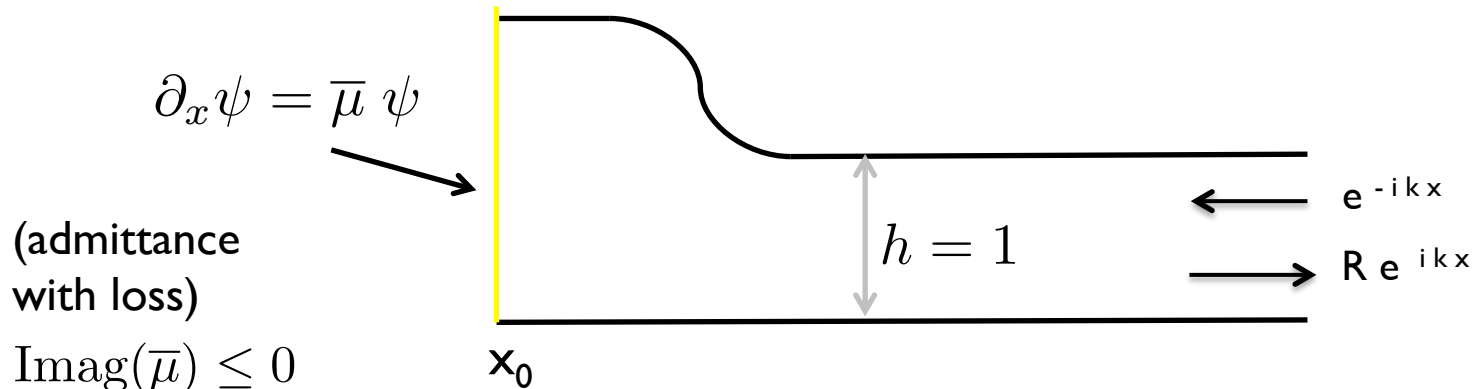
Solution with **total absorption**

Total absorption

In the following $k < \pi$: only the mode 0 is propagating

Absorption at **low frequencies**

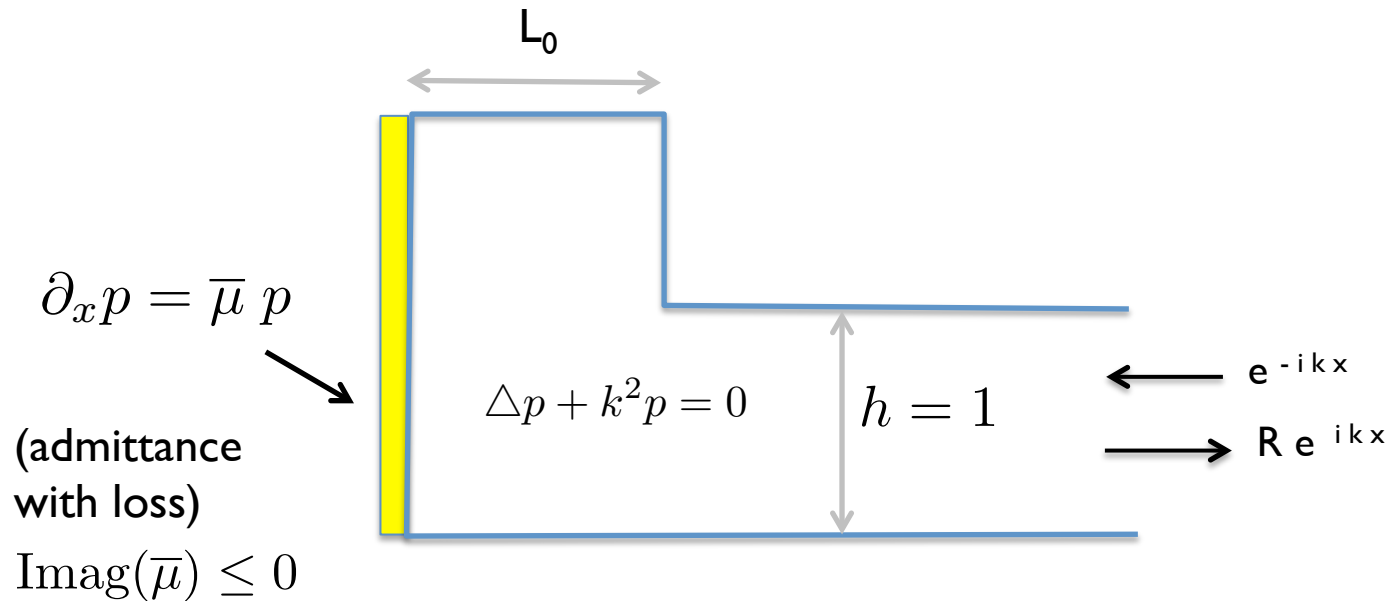
$$Y_r = \begin{pmatrix} ik & & & \\ & -|k_1| & & \\ & & -|k_2| & \\ & & & \dots \end{pmatrix}$$



To obtain **total absorption** $R=0$:

choose the admittance μ as the eigenvalue problem of $Y(x_0)$ (outgoing wave pb)

Results



Where μ is one of the $\mu_0, \mu_1, \mu_2, \dots$ (eigenvalues of $Y(x_0)$ for outgoing waves)

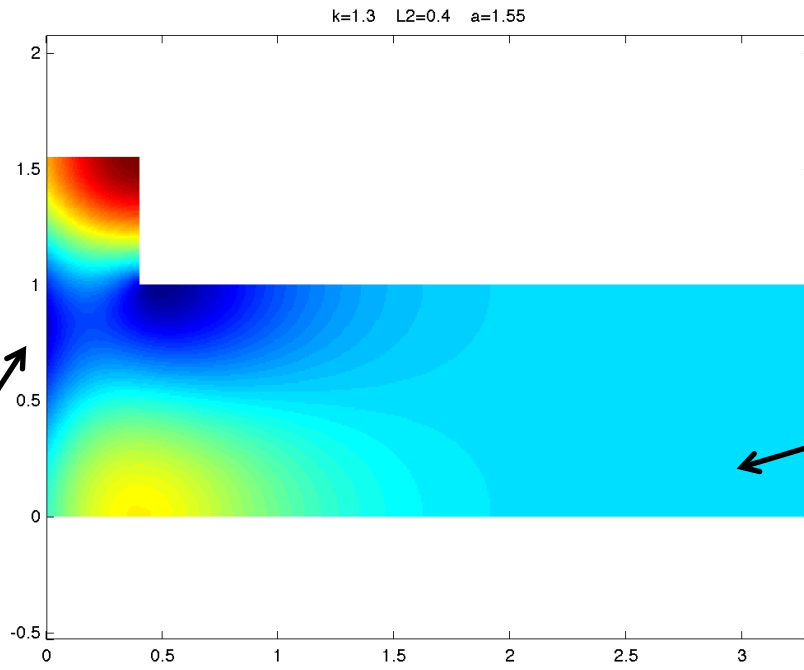
Remark :

For a straight strip, $Y(x_0) = Y_r$ and only $\mu_0 = ik$ gives an absorbing admittance.

This value (ik) is too large for thin coating of porous media at low freqs:

thus, we would like $|\mu| < k$

Total absorption example



$|p(x,y)|$

Mode matching

$|p(x,y)| = 1$
(constant since no reflection)

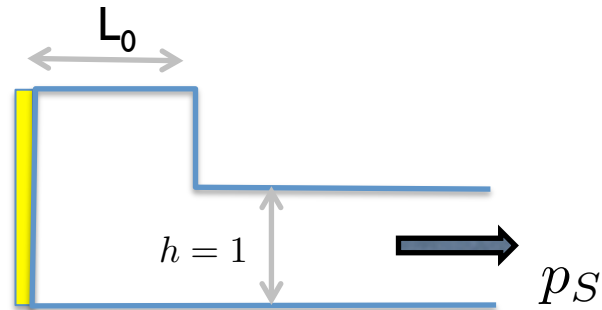
$$\partial_x p = \bar{\mu} p$$



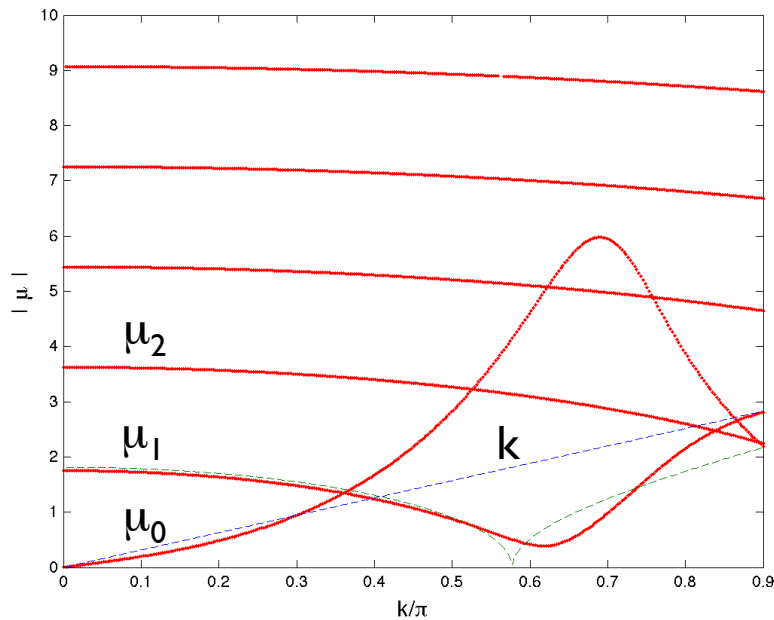
eigenvalues μ_n as a function of frequency k

$$L_0 = 1$$

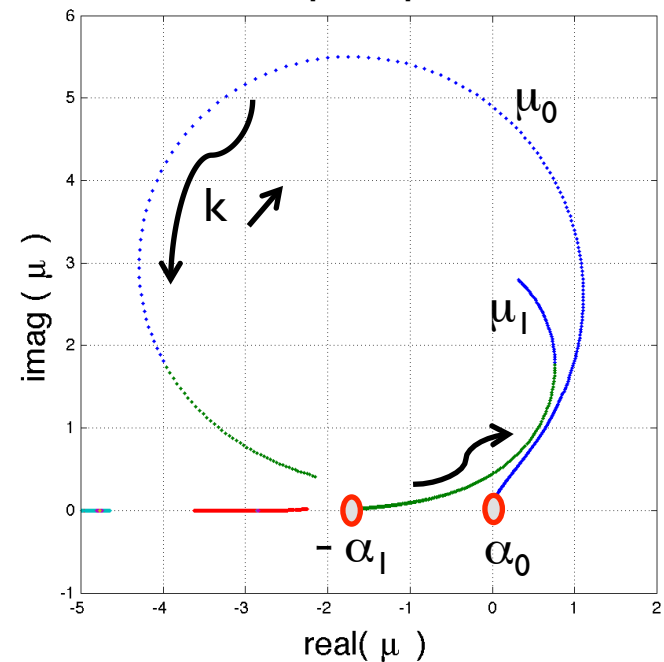
$$\partial_x p = \mu p$$



Absolute value of μ vs k

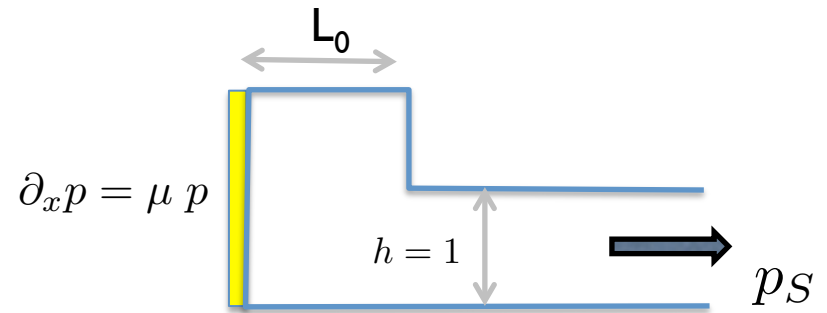


Complex plane of μ

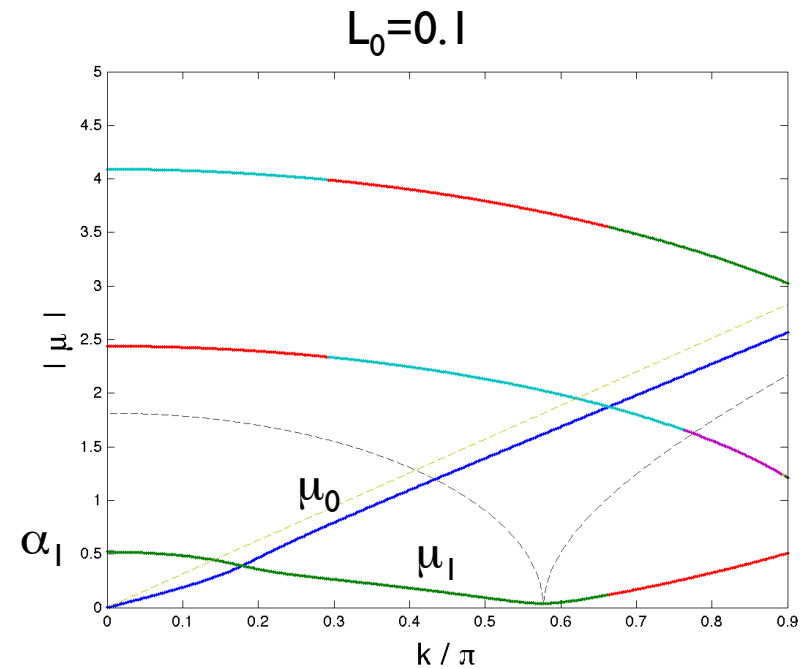
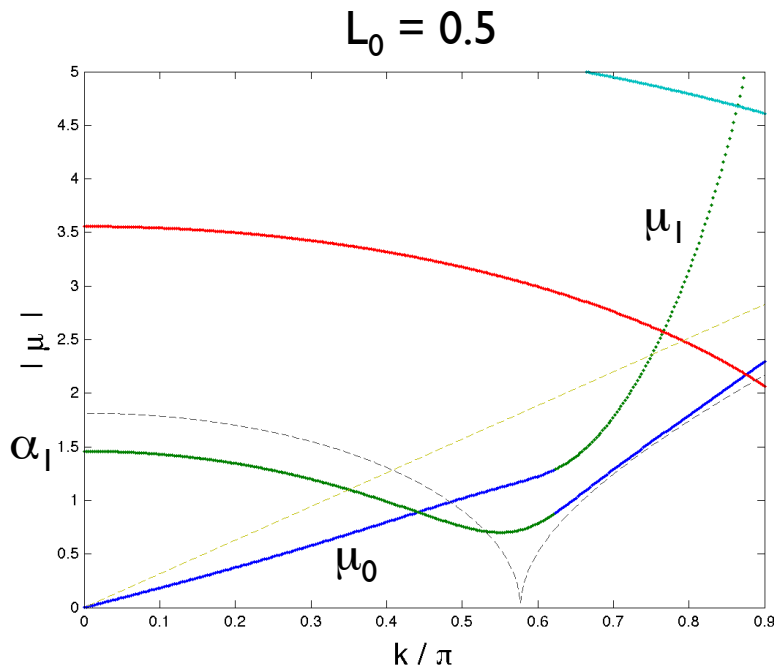


$-\alpha_n$: value of μ_n for $k \rightarrow 0$ ($\alpha_0 = 0$ and $\alpha_n > 0$)

Diminishing the length cavity L_0 : α_1 decreases

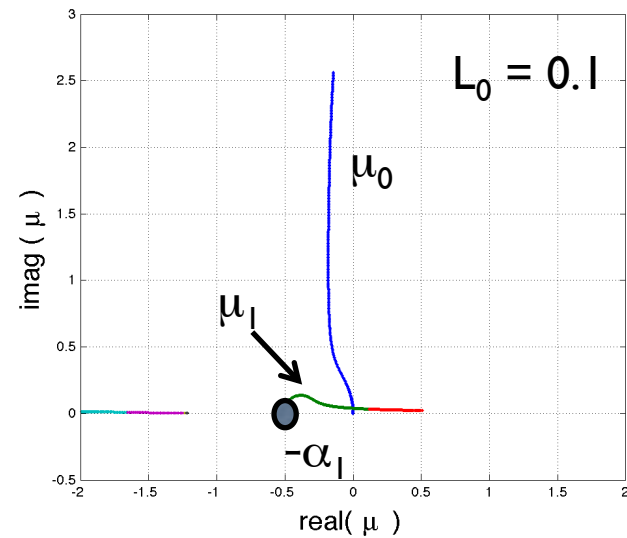
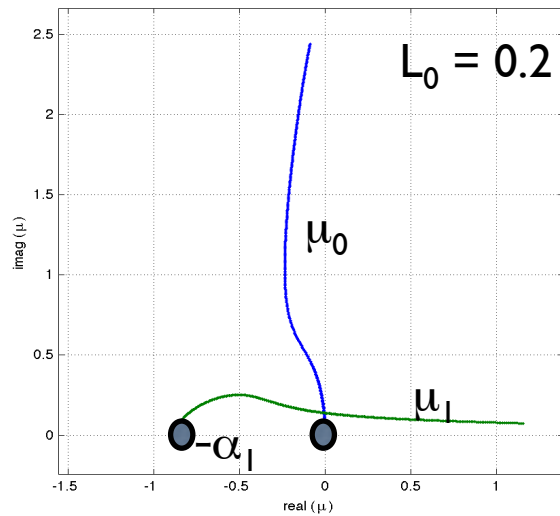
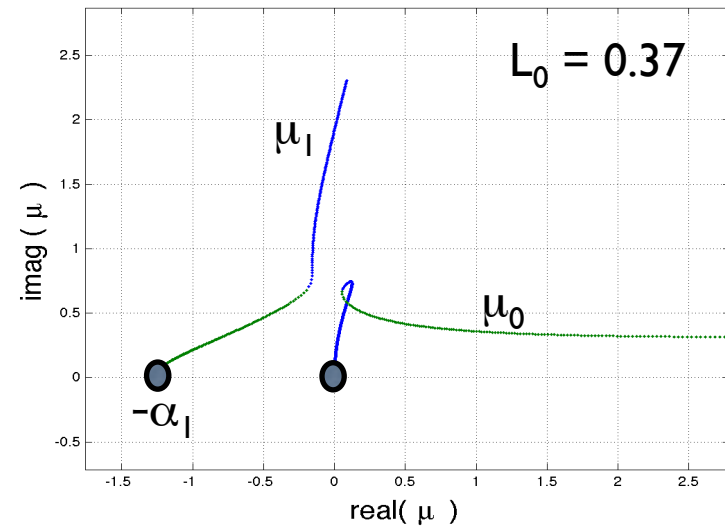
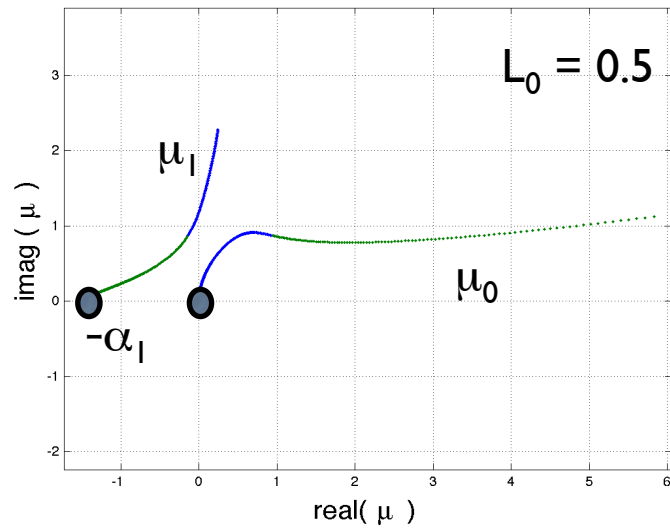


Absolute values of μ_n vs frequency k

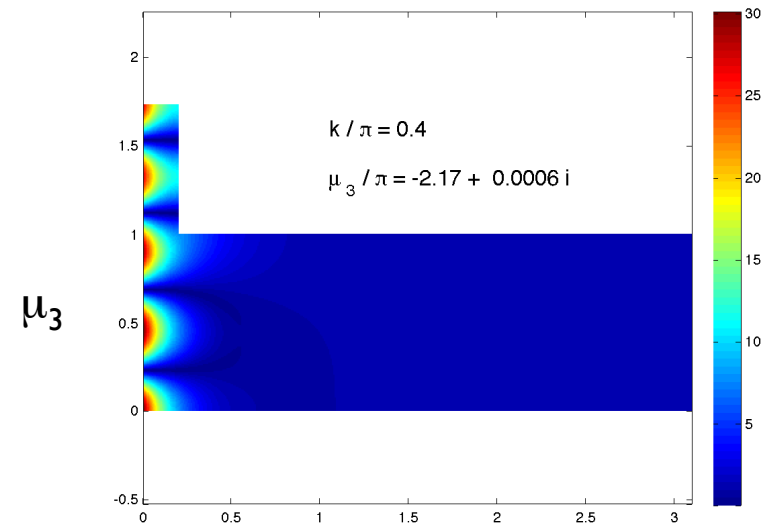
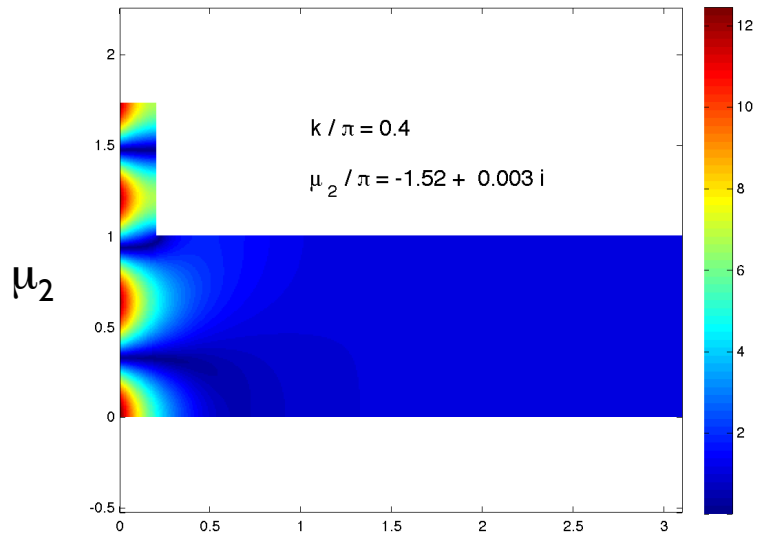
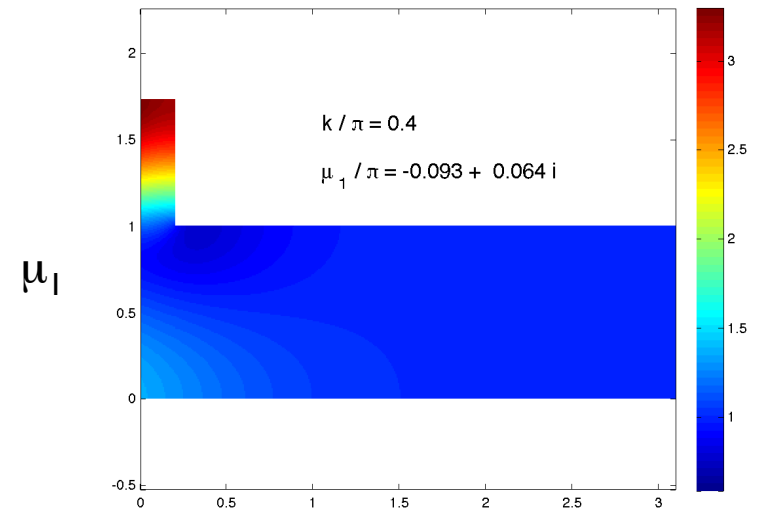
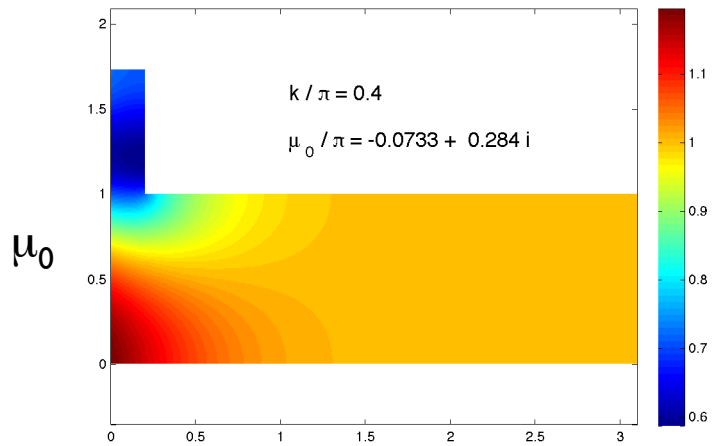


Possible : μ_n smaller than k

eigenvalues μ_n as a function of frequency k : complex μ plane



Different μ_n : different admittance BC for total absorption (for the same k)

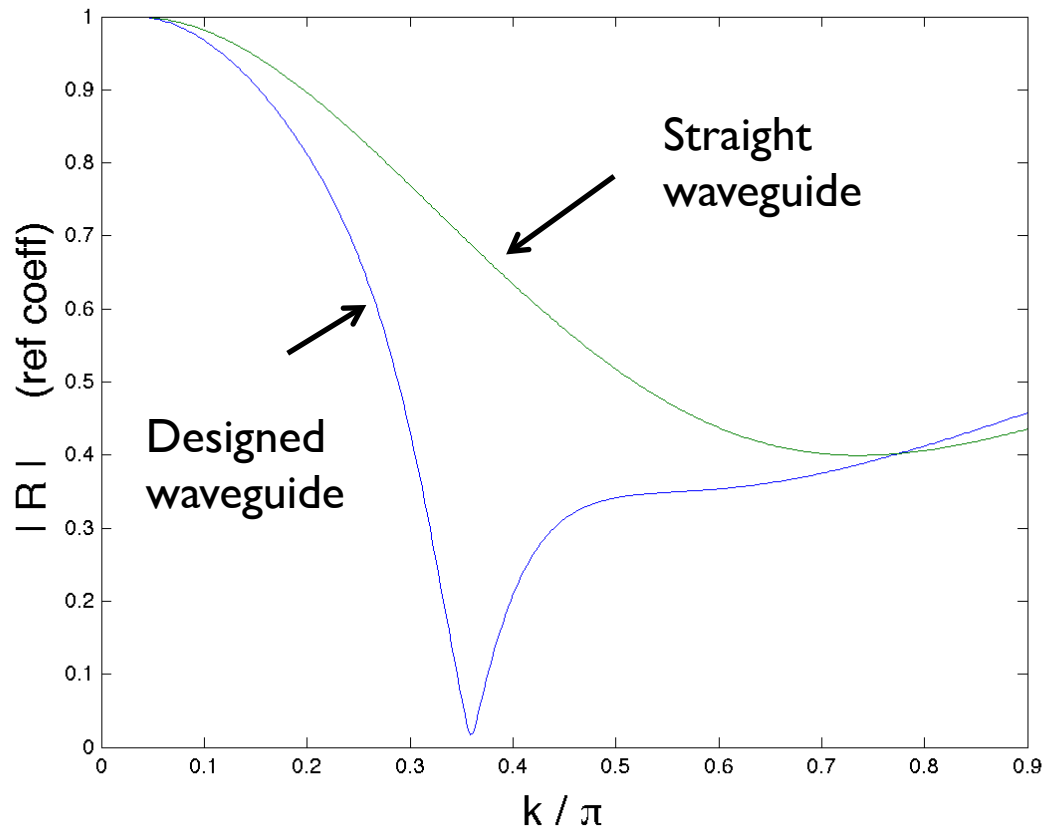


Reflection coefficient vs frequency k

same admittance

Comparison between

- Straight waveguide
- Heterogeneous waveguide designed to totally absorb at $k=0.37\pi$ (admittance = thin porous coating similar to Delany-Basley)



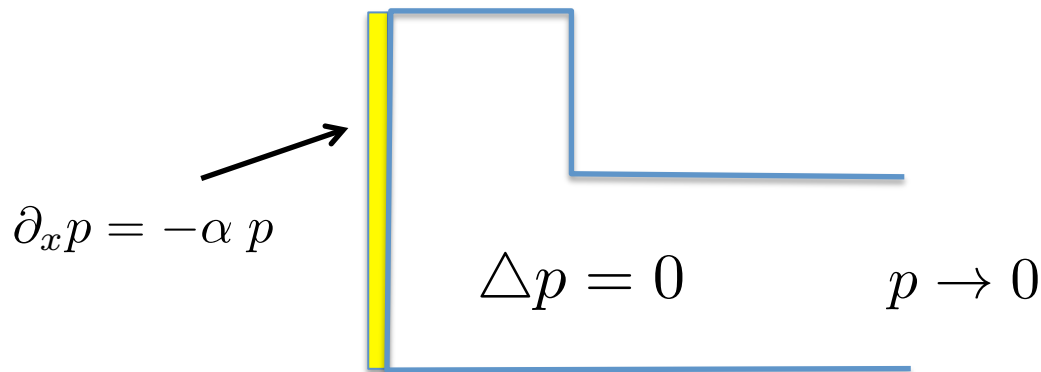
Connection with water waves

To get low values of μ_n at low frequency :

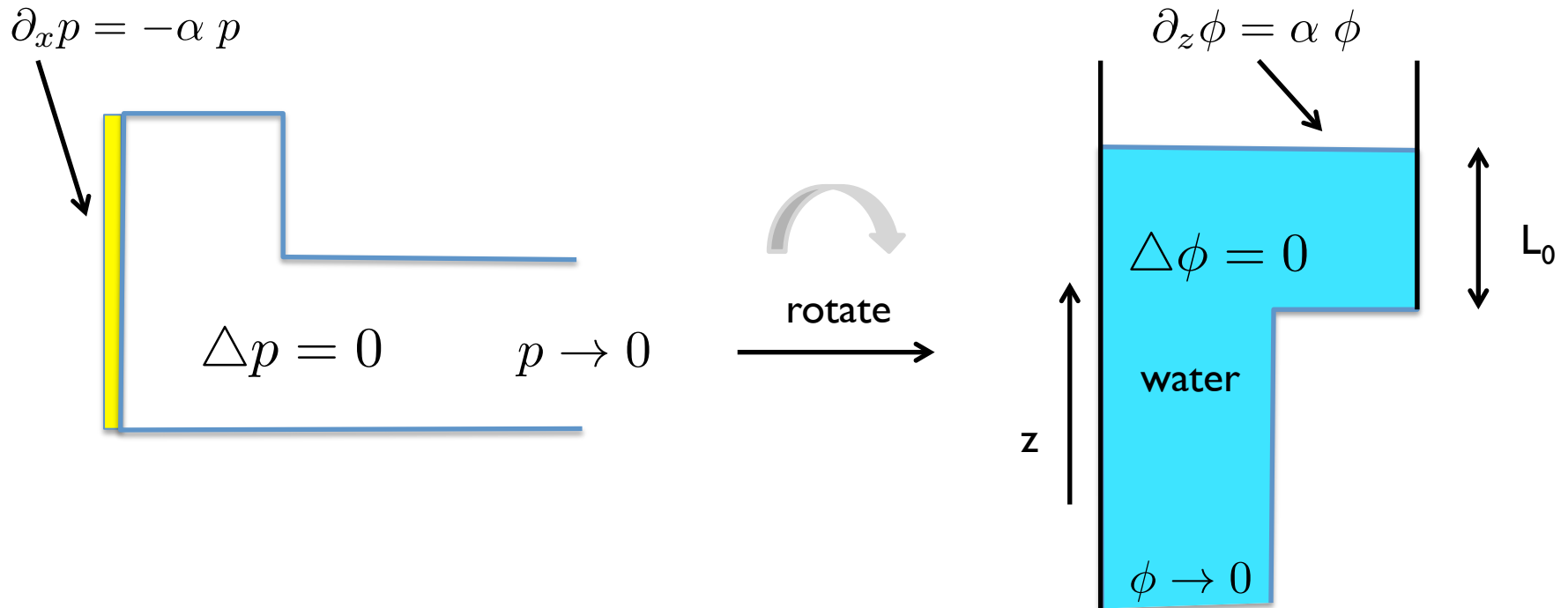
$\alpha_n = \lim \mu_n$ as $k \rightarrow 0$ is important

When α_n is small , total absorption by “small admittance”
is possible at low frequency (i.e. absorption by thin structure)

Asymptotic expansion ($k \rightarrow 0$) : problem for α_n



Connection with water waves



α is the eigenvalue of the **sloshing for water waves in cavity**

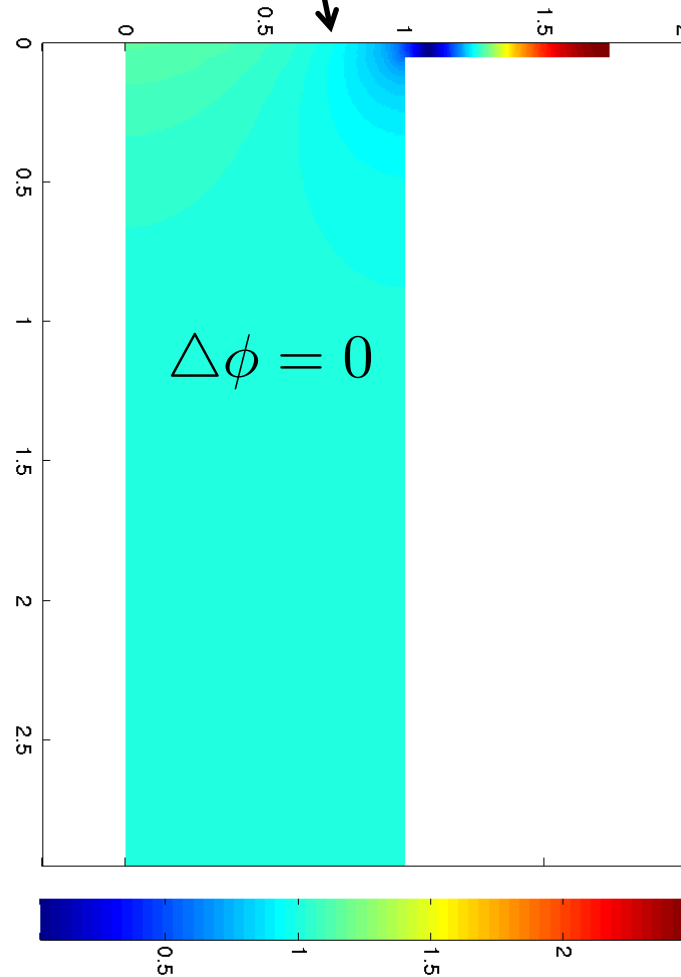
When $L_0 \rightarrow 0$: $\alpha_1 \rightarrow 0$ (shallow water has small wave speed)

$$\alpha_1 \simeq L_0 (\pi/d)^2$$

When $L_0 \rightarrow 0$: $\alpha_1 \rightarrow 0$ (shallow water has small wave speed)

$$\alpha_1 \simeq L_0 (\pi/d)^2$$

$$\partial_z \phi = \alpha \phi$$



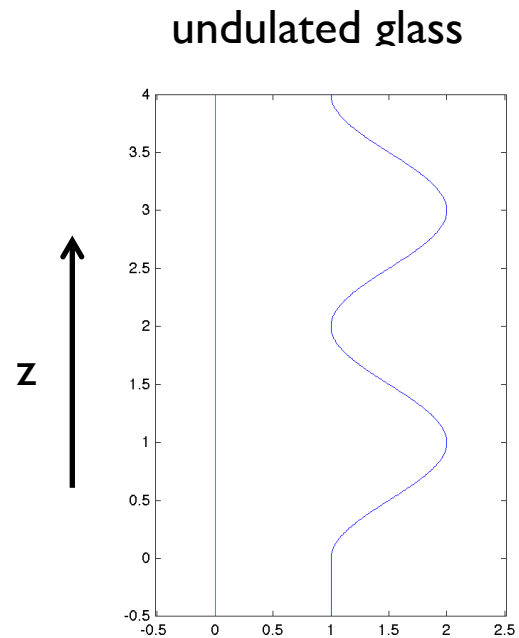
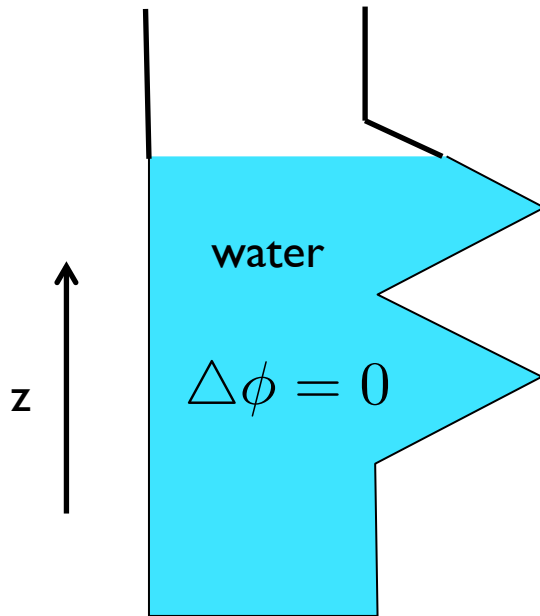
“flute mode”

Riccati for sloshing

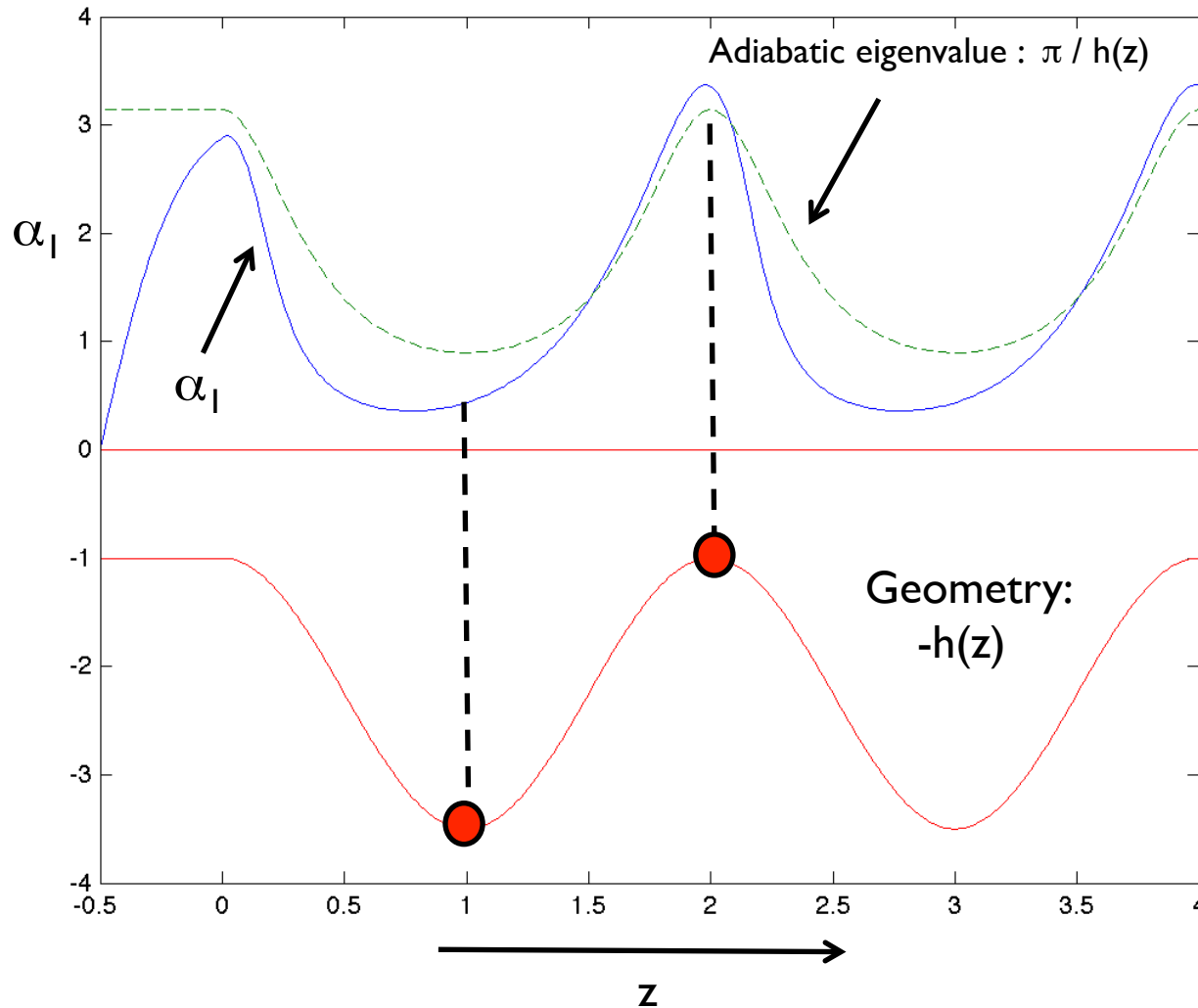
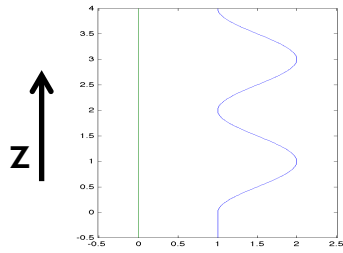
Integrating Riccati with z gives
the sloshing frequency at each z
(eigenvalues of $Y(z)$)

$$Y' = -\cancel{K^2} - Y^2 + YF + F^T Y$$

Riccati equation



Shloshing eigenvalue α_1 as a function of z



Domain Monotonicity
for sloshing
at red points ●

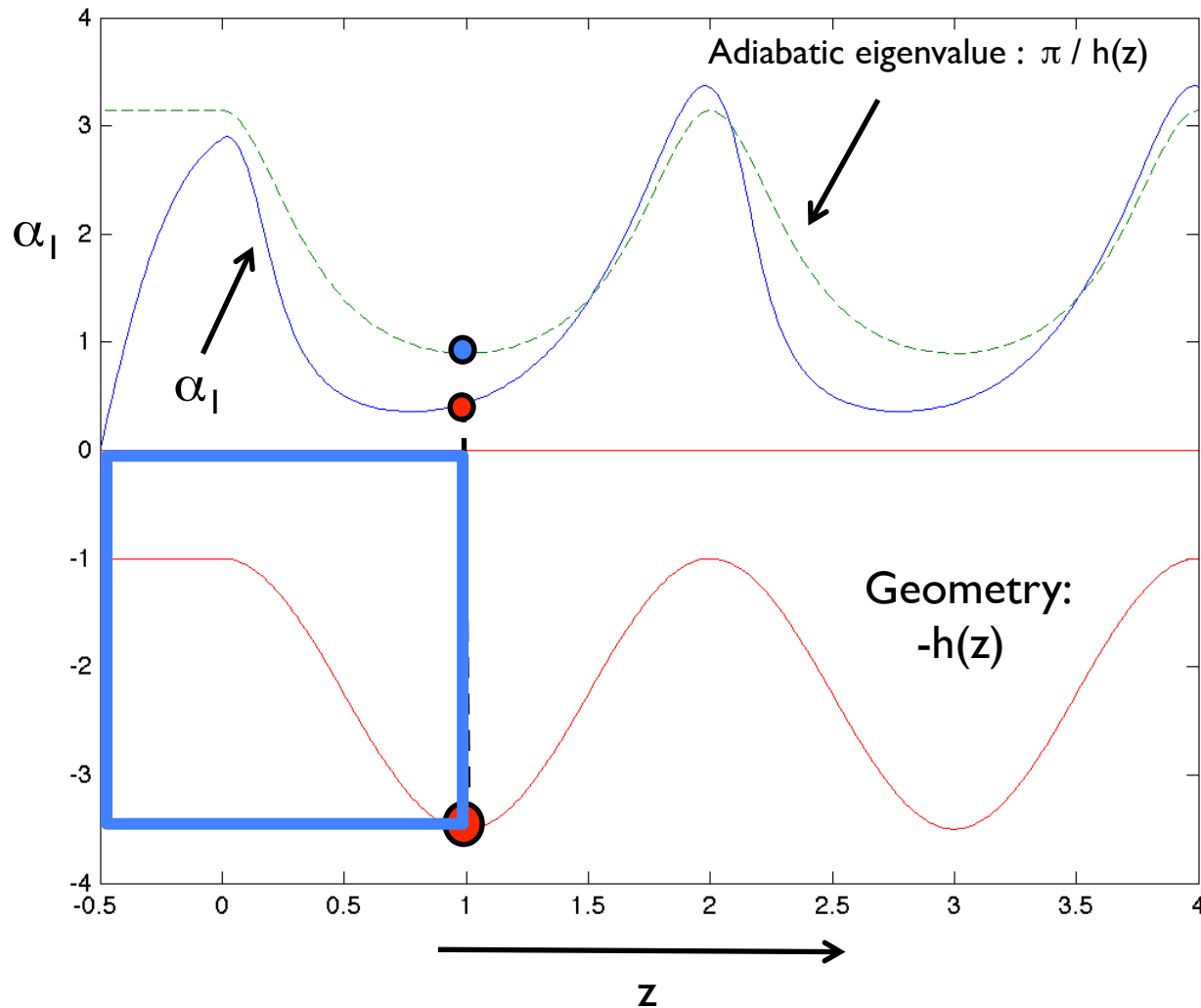
$$\text{If } \hat{\Omega} \subset \Omega$$

$$\text{then } \hat{\alpha}_1 \leq \alpha_1$$

(Ω volume of the cavity)

π/h : eigenvalue of
rectangular cavity

$$\alpha_1 \leq \pi/h$$



Domain Monotonicity
for sloshing
at red points ●

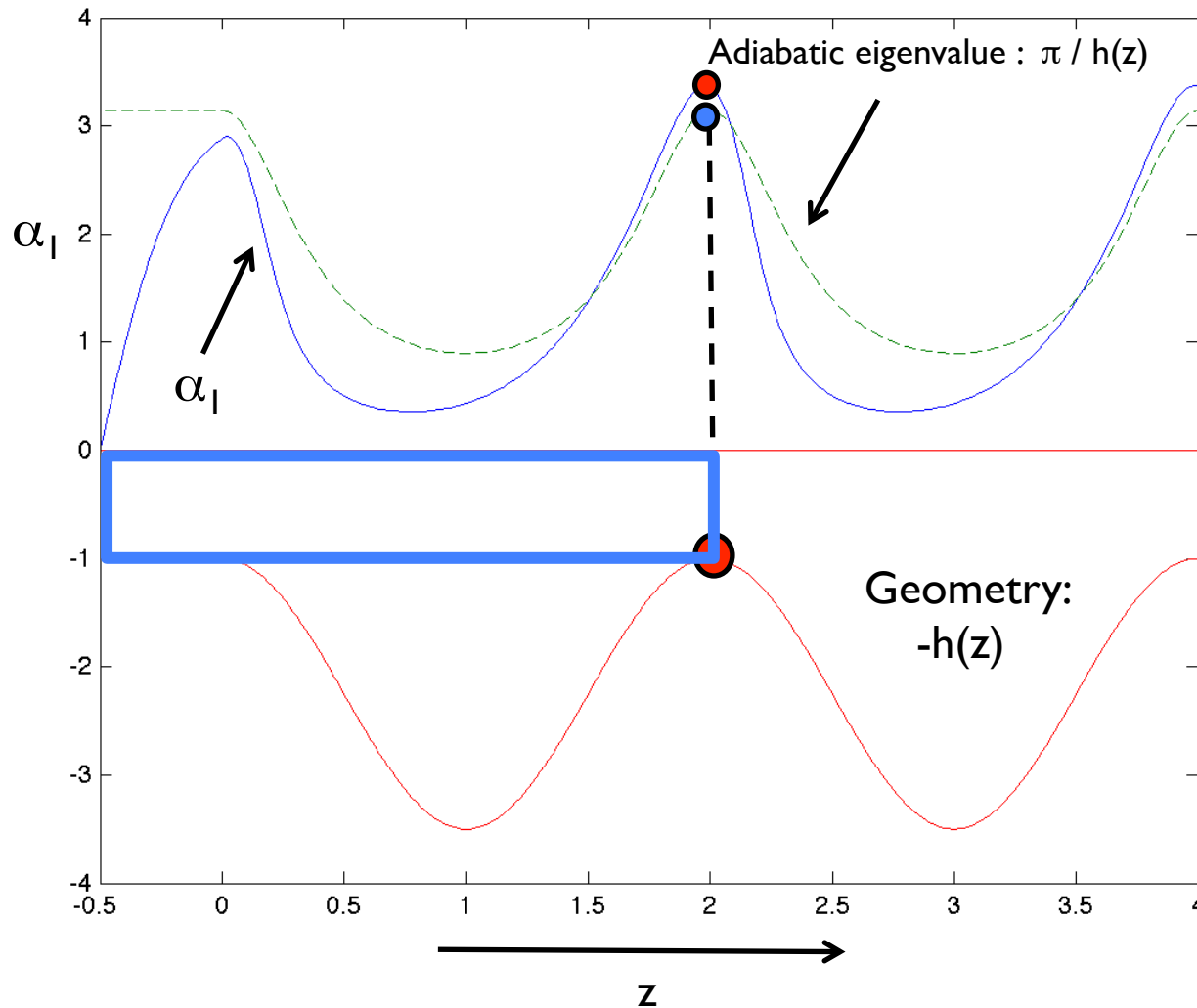
If $\hat{\Omega} \subset \Omega$

then $\hat{\alpha}_1 \leq \alpha_1$

(Ω volume of the cavity)

π/h : eigenvalue of
rectangular cavity

$$\pi/h \leq \alpha_1$$



Domain Monotonicity
for sloshing
at red points ●

If $\hat{\Omega} \subset \Omega$

then $\hat{\alpha}_1 \leq \alpha_1$

(Ω volume of the cavity)

π/h : eigenvalue of
rectangular cavity

Links to high spot (hot spot) problems:

Importance of the angle between the free surface and the wall

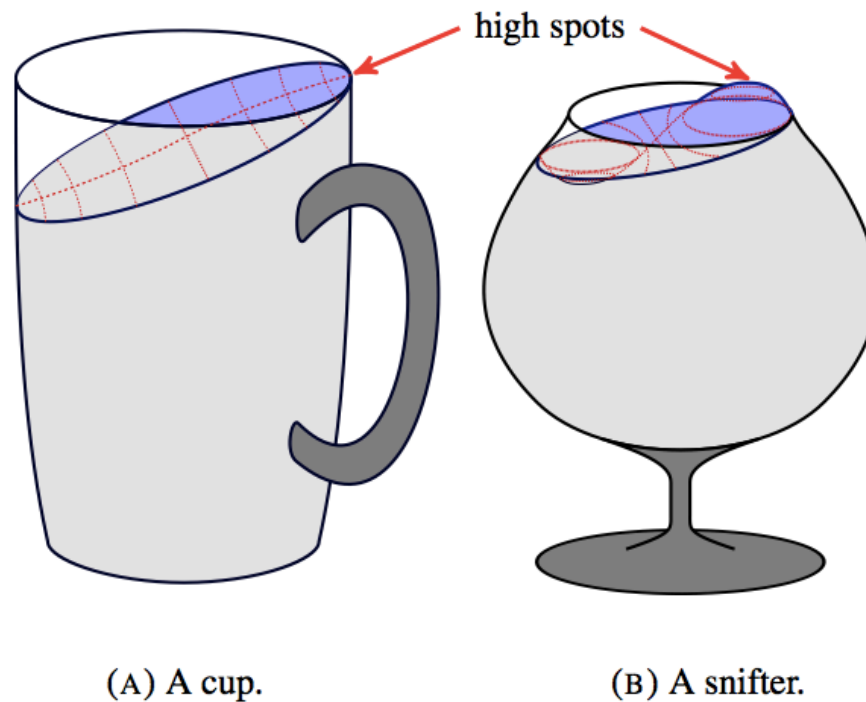


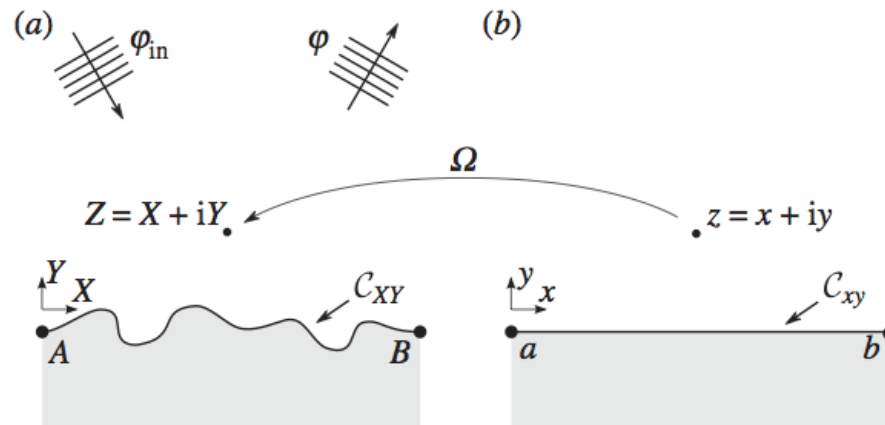
FIGURE 1. High spots in a coffee cup and a snifter.

Kulczycki, T., Kwasnicki, M., & Siudeja, B. (2013).

Spilling from a cognac glass. *arXiv preprint arXiv:1311.7296*.

Remark (rough surface)

The same (multimodal Riccati) can be done for rough surface by conformal mapping

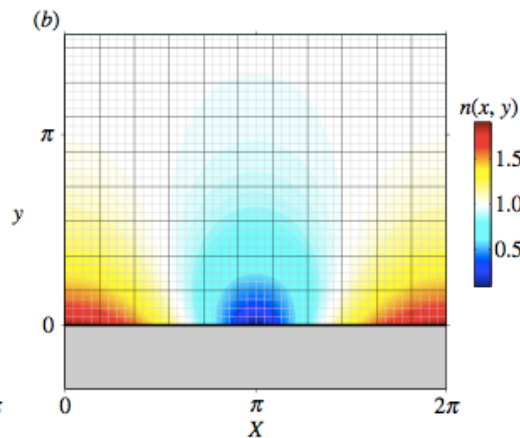
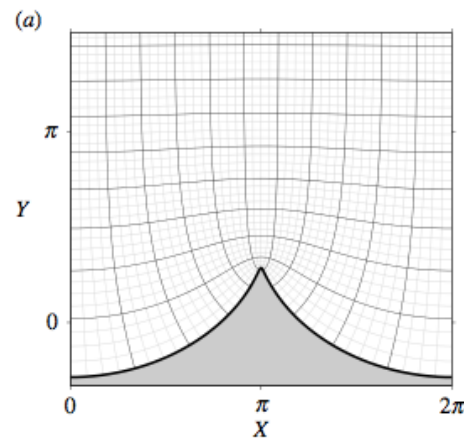


$$\Delta_{XY}\Phi + k^2\Phi = 0$$

$$\Delta_{xy}\Phi + k^2 n^2 \Phi = 0$$

“Refractive index”

$$n(x, y) = \left| \frac{dZ}{dz} \right|$$

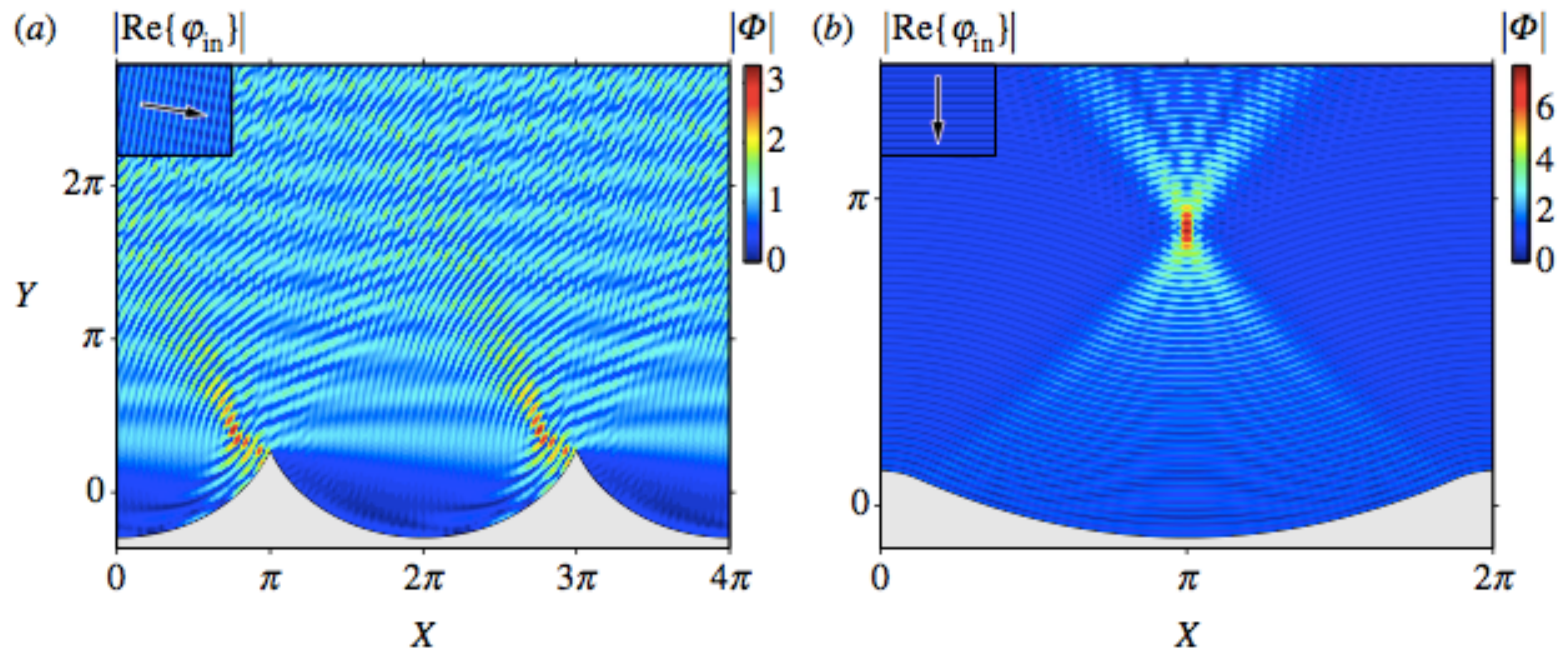


Riccati eq. for the
Admittance matrix
(DtN operator)

$$dY/dy = f(Y, y)$$

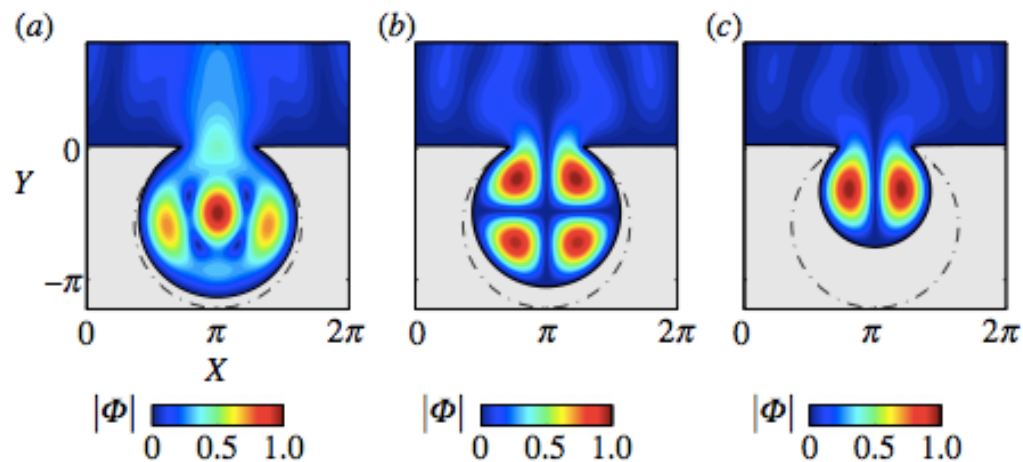
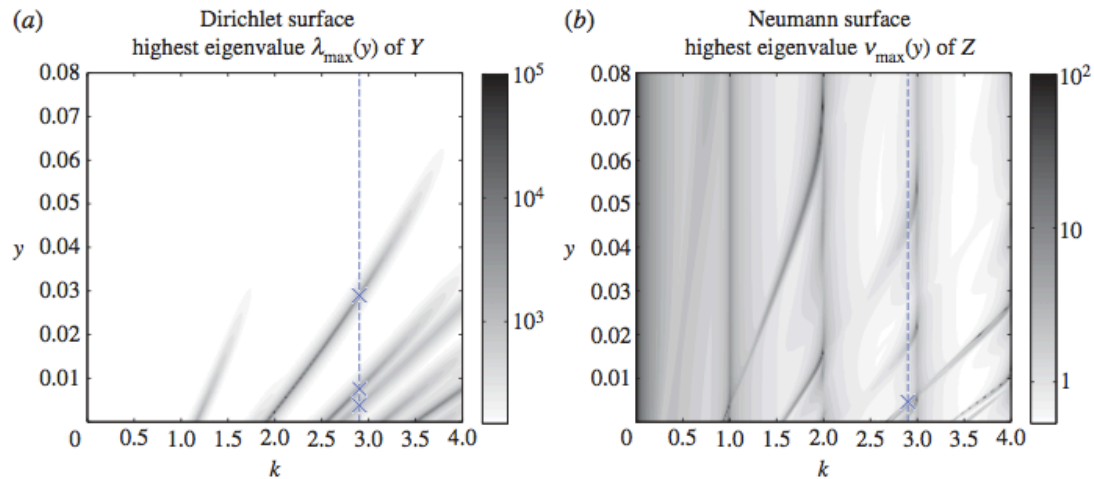
Remark (rough surface)

Multimodal Riccati to compute the scattering by the surface



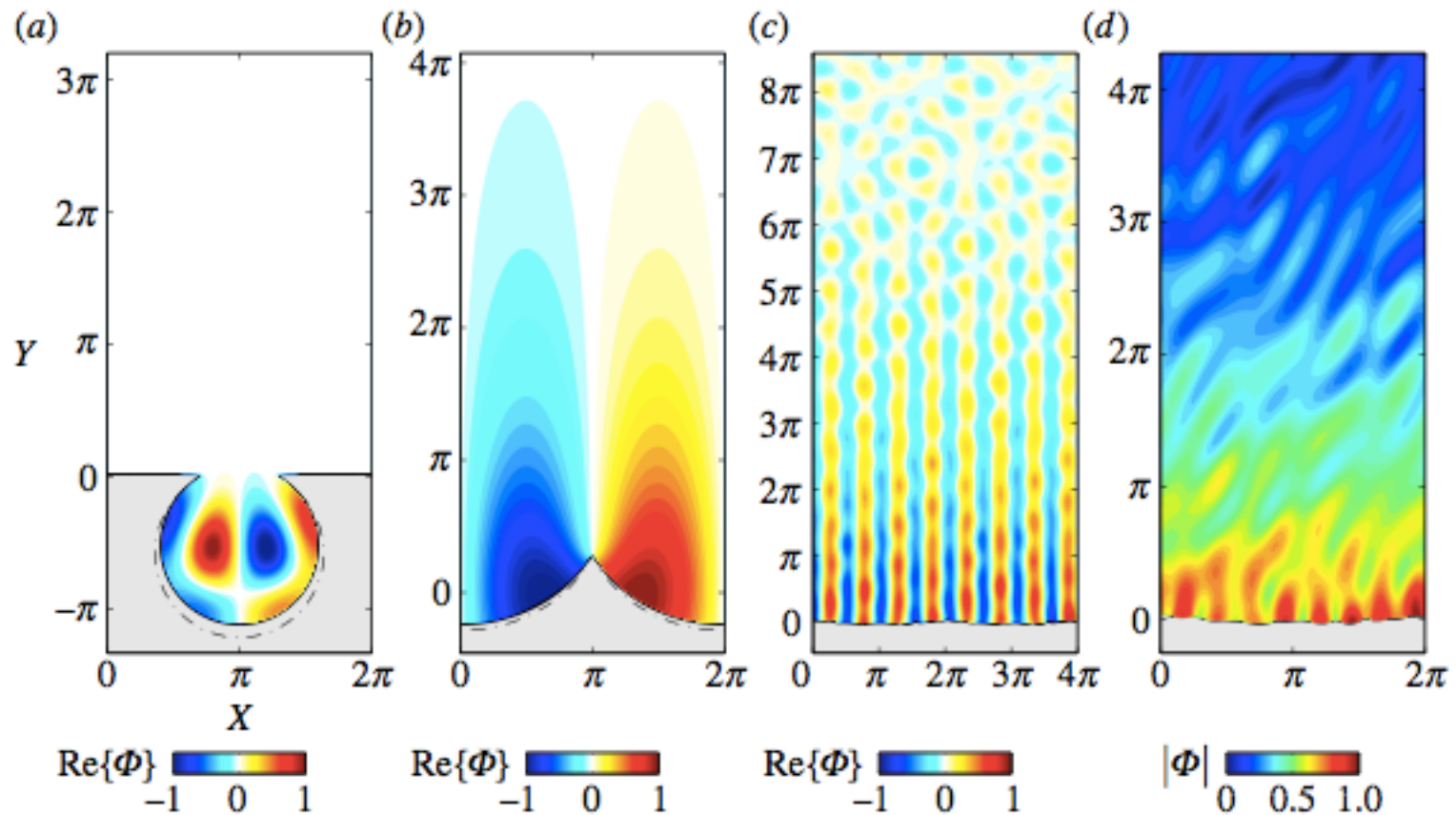
Remark (rough surface)

From Riccati : Eigenvalues of $Y(y)$ can give total absorption or trapped surface waves



Remark (rough surface)

Exemples of surface waves “trapped by the roughness”
(also called BIC)



V. Pagneux, Multimodal admittance method and singularity behaviour at high frequencies, *J. Comp. App. Maths.* **234** (6), 1834-1841 (2010)